

working with Keith on  
***Gravitino Dark Matter***  
models

Vassilis C. Spanos  
UoA

Olivefest 2017, Mnpls 19/5/2017

# Outline

- Keith's early Gravitino days
- Now: why still SUSY, Supersymmetric DM scenarios
- Direct, indirect searches of DM
- Gravitino Dark Matter
- Gravitino thermal production and reheating temperature
- Recap

## SUPERSYMMETRIC RELICS FROM THE BIG BANG\*

John ELLIS and J. S. HAGELIN

*Stanford Linear Accelerator Center, Stanford University, Stanford, California 94305, USA*

D. V. NANOPOULOS, K. OLIVE<sup>†</sup>, and M. SREDNICKI<sup>‡</sup>

*CERN, CH-1211 Geneva 23, Switzerland*

Received 16 September 1983

(Revised 15 December 1983)

We consider the cosmological constraints on supersymmetric theories with a new, stable particle. Circumstantial evidence points to a neutral gauge/Higgs fermion as the best candidate for this particle, and we derive bounds on the parameters in the lagrangian which govern its mass and couplings. One favored possibility is that the lightest neutral supersymmetric particle is predominantly a photino  $\tilde{\gamma}$  with mass above  $\frac{1}{2}$  GeV, while another is that the lightest neutral supersymmetric particle is a Higgs fermion with mass above 5 GeV or less than  $O(100)$  eV. We also point out that a gravitino mass of 10 to 100 GeV implies that the temperature after completion of an inflationary phase cannot be above  $10^{14}$  GeV, and probably not above  $3 \times 10^{12}$  GeV. This imposes constraints on mechanisms for generating the baryon number of the universe.

Nuclear Physics B255 (1985) 495–504  
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## GRAVITINOS AS THE COLD DARK MATTER IN AN $\Omega = 1$ UNIVERSE

Keith A. OLIVE and David N. SCHRAMM<sup>1</sup>

*Astrophysics Theory Group, Fermilab, PO Box 500, Batavia, IL 60510, and <sup>1</sup> Departments of Astronomy and Astrophysics, Physics and EFI, The University of Chicago, IL 60637, USA*

Mark SREDNICKI

*Department of Physics, The University of California, Santa Barbara, CA 93106, USA*

Received 26 October 1984

In attempts to simultaneously have galaxy formation with adiabatic fluctuations, dark galactic halos and a critical  $\Omega = 1$ , recent models have been proposed with a heavy (cold) species which decays and a lighter (hot or warm) one which is stable. Surprisingly enough, these models are very constraining. Independent of the particle physics model, the decay width of the heavy particle must be  $O(10^{-40})$  GeV. Such a scale is only natural in theories involving gravity where one expects  $\Gamma \sim M_H^3 / M_P^2$ , where  $M_H$  is the mass of the heavy particle and  $M_P$  is the Planck mass. In this paper, we suggest that the heavy particle might be the gravitino which is present in all supergravity theories. This model would then require that the gravitino decay products not include photons, indicating that the lightest supersymmetric particle (LSP) must be something other than the photino. An acceptable candidate for the LSP might be the axino, the supersymmetric partner of the axion.

Volume 145B, number 3,4

PHYSICS LETTERS

20 September 1984

## COSMOLOGICAL GRAVITINO REGENERATION AND DECAY

John ELLIS, Jihn E. KIM<sup>1</sup> and D.V. NANOPOULOS

*CERN, Geneva, Switzerland*

Received 29 May 1984

We present a detailed study of gravitino production and decay subsequent to cosmological inflation. We calculate the cross sections for gravitino production in collisions of particles in the supersymmetric standard model, and use them to calculate the regenerated abundance of gravitinos as a function of the maximum reheating temperature  $T_{\max}$ . An upper limit on the gravitino mass density during cosmological nucleosynthesis requires  $T_{\max} < 0.90 \times 10^{16}$  GeV and considerations of the entropy released when gravitinos decay require  $T_{\max} < 2.2 \times 10^{13}$  GeV, while more careful analyses of their decay products' disruptive effects on light nuclei and on the microwave background radiation suggest  $T_{\max} < 10^9 - 10^{10}$  GeV.

VOLUME 55, NUMBER 21

PHYSICAL REVIEW LETTERS

18 NOVEMBER 1985

## Diffuse Cosmic Gamma-Ray Background as a Probe of Cosmological Gravitino Regeneration and Decay

Keith A. Olive

*Astrophysics Theory Group, Fermilab, Batavia, Illinois 60510*

and

Joseph Silk

*Department of Astronomy, University of California, Berkeley, California 94720*  
(Received 22 March 1985)

We predict the presence of a spectral feature in the isotropic cosmic gamma-ray background associated with gravitino decays at high red shifts. With a gravitino abundance that falls in the relatively narrow range expected for thermally regenerated gravitinos following an inflationary epoch in the very early universe, gravitinos of mass several gigaelectronvolts are found to yield an appreciable flux of 1–10-MeV diffuse gamma rays.

PACS numbers: 98.70.Vc, 11.30.Pb, 14.80.Pb, 95.30.Cq



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Physics Letters B 588 (2004) 7–16

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PHYSICS LETTERS B

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[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)

## Gravitino dark matter in the CMSSM

John Ellis <sup>a</sup>, Keith A. Olive <sup>b</sup>, Yudi Santoso <sup>b</sup>, Vassilis C. Spanos <sup>b</sup>

<sup>a</sup> *TH Division, CERN, Geneva, Switzerland*

<sup>b</sup> *William I. Fine Theoretical Physics Institute, University of Minnesota, Minneapolis, MN 55455, USA*

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### Abstract

We consider the possibility that the gravitino might be the lightest supersymmetric particle (LSP) in the constrained minimal extension of the Standard Model (CMSSM). In this case, the next-to-lightest supersymmetric particle (NSP) would be unstable, with an abundance constrained by the concordance between the observed light-element abundances and those calculated on the basis of the baryon-to-entropy ratio determined using CMB data. We modify and extend previous CMSSM relic neutralino calculations to evaluate the NSP density, also in the case that the NSP is the lighter stau, and show that the constraint from late NSP decays is respected only in a limited region of the CMSSM parameter space. In this region, gravitinos might constitute the dark matter.

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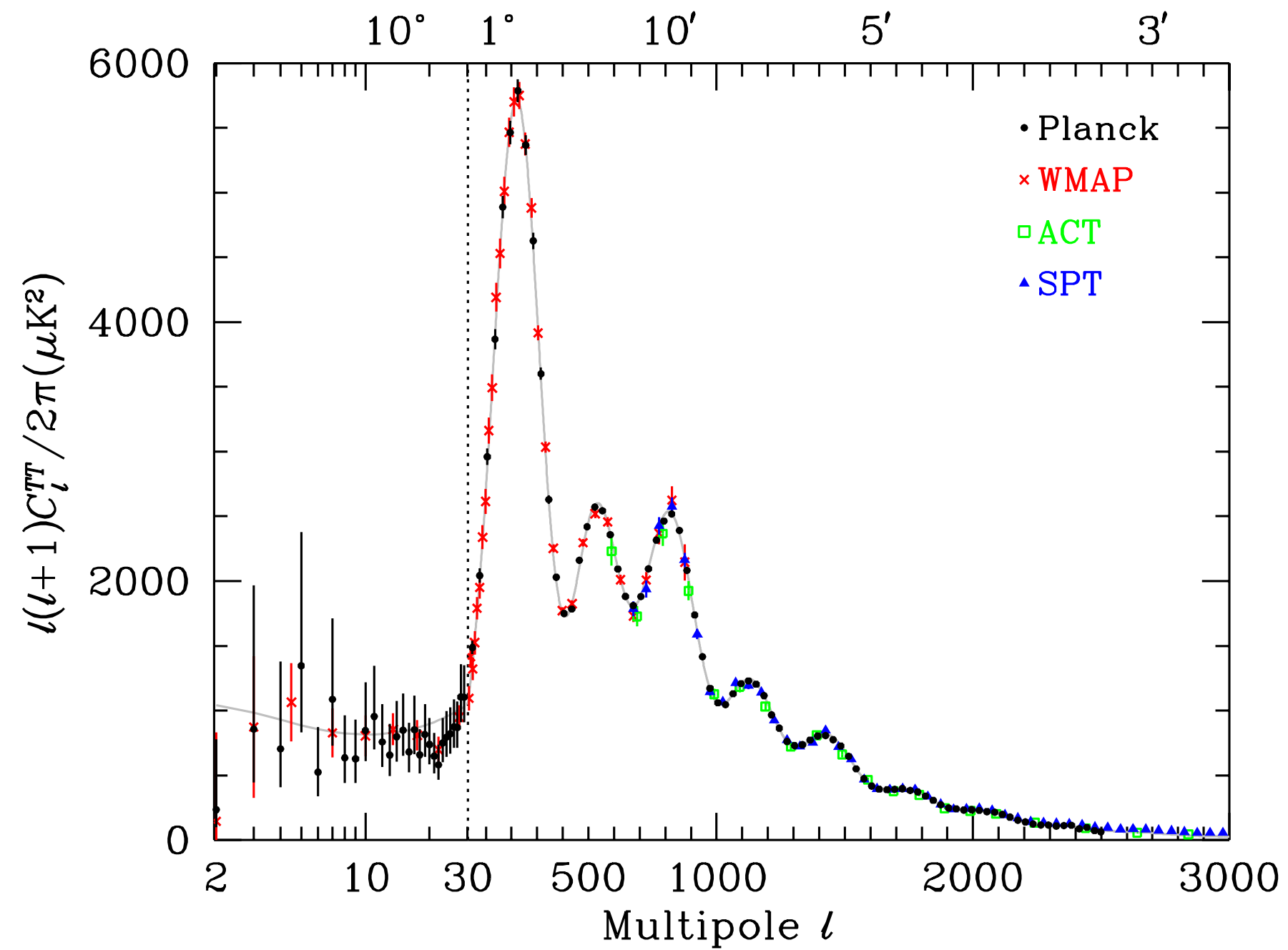
# SUSY and Higgs mass

$$m_h = \sqrt{\frac{1}{2}[M_A^2 + M_Z^2 - \sqrt{(M_A^2 + M_Z^2)^2 - 4M_A^2 M_Z^2 \cos^2 \beta}]} < M_Z$$

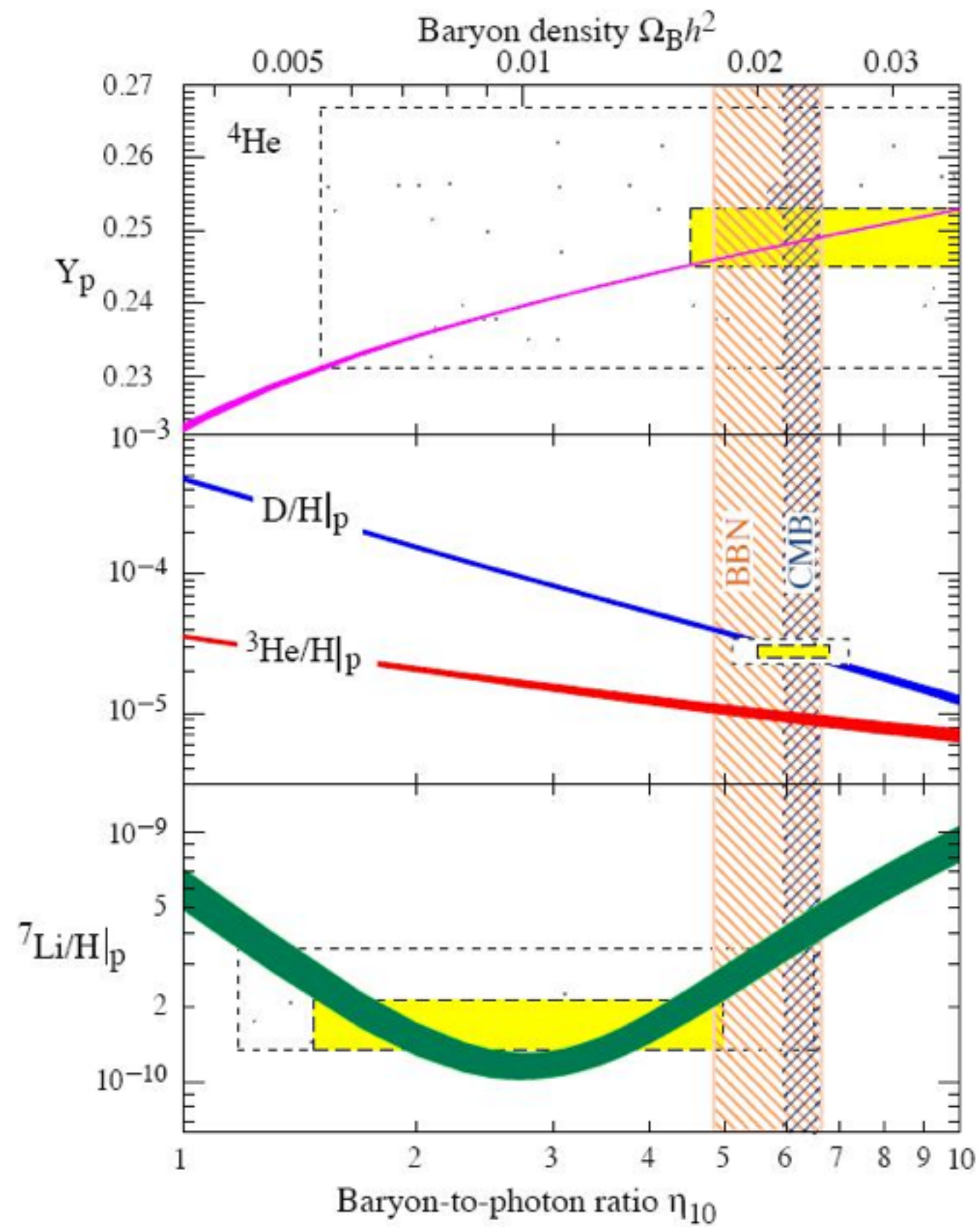
$$\tan \beta = \frac{v_2}{v_1}$$

$$\begin{aligned} \Delta m_h^2 \simeq & \frac{3}{4\pi^2} \frac{m_t^4}{v^2} \left[ \log \left( \frac{m_{\tilde{t}}^2}{m_t^2} \right) + \frac{X_t^2}{m_{\tilde{t}}^2} - \frac{X_t^4}{12m_{\tilde{t}}^4} \right] \\ & - \frac{3}{48\pi^2} \frac{m_b^4}{v^2} \frac{t_\beta^4}{(1 + \epsilon_b t_\beta)^4} \frac{\mu^4}{m_{\tilde{b}}^4} \\ & - \frac{1}{48\pi^2} \frac{m_\tau^4}{v^2} \frac{t_\beta^4}{(1 + \epsilon_\ell t_\beta)^4} \frac{\mu^4}{m_{\tilde{\tau}}^4} \end{aligned} \quad X_t = A_t + \mu / \tan \beta \approx A_t$$

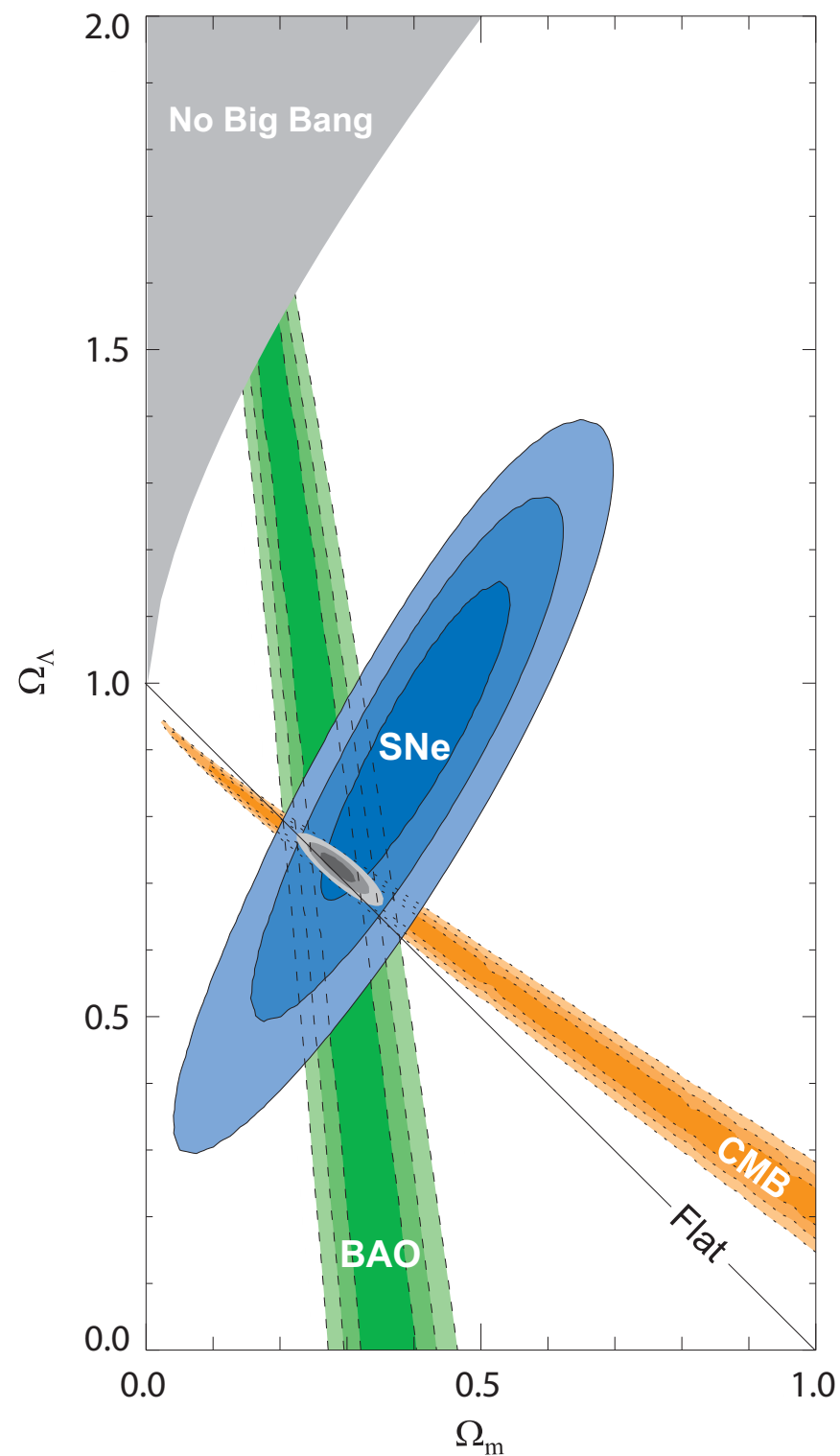
# DM/ Evidence



## DM/ Evidence



## DM/ Evidence

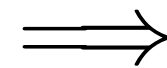


$$\Omega_{\text{tot}} \sim 1.0$$

$$\Omega_{\text{M}} = 0.315 \pm 0.016$$

$$\Omega_{\Lambda} = 0.685 \pm 0.017$$

$$\Omega_{\text{B}} = 0.0499 \pm 0.0022$$

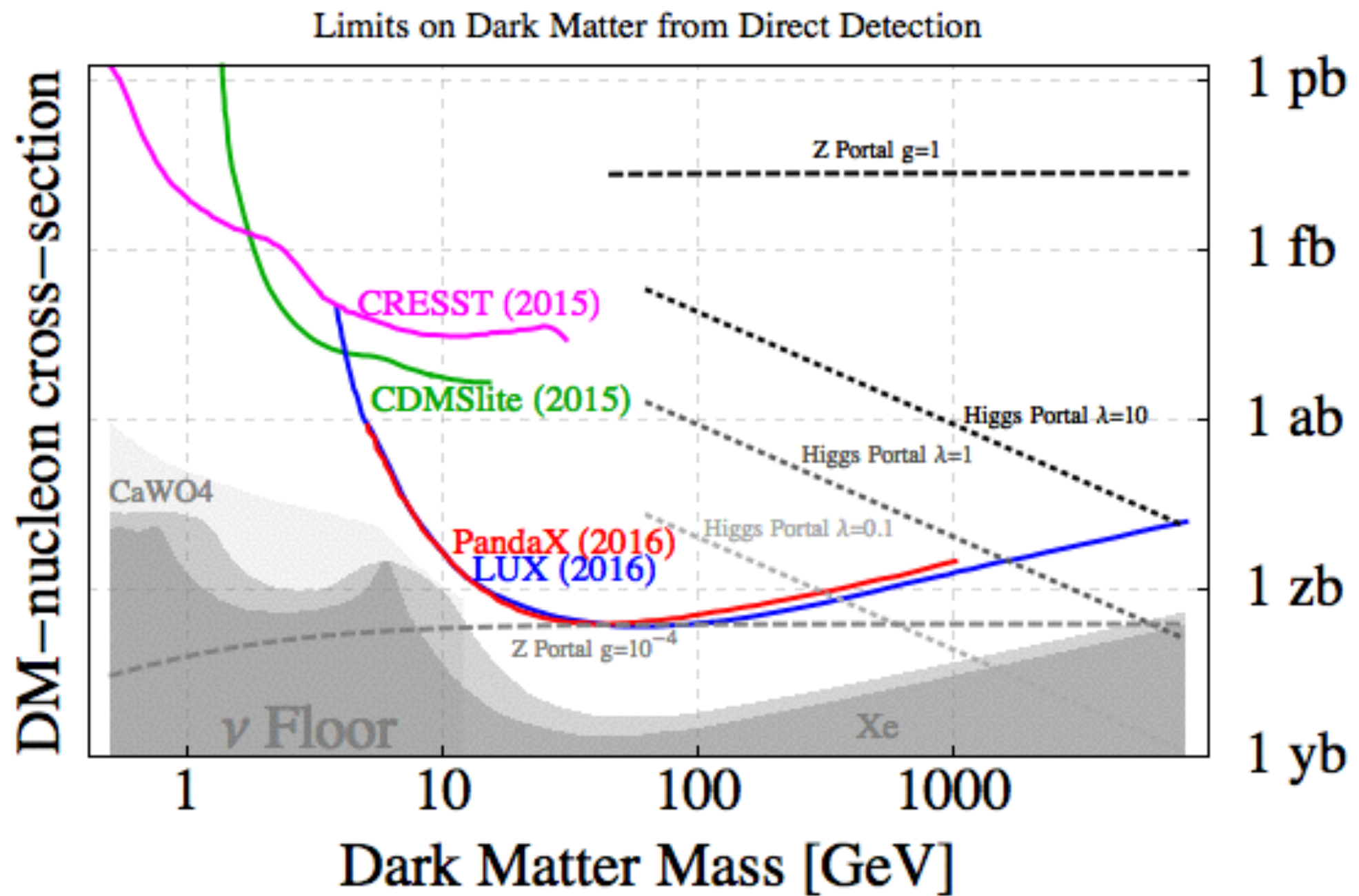


$$\Omega_{\text{DM}} = \Omega_{\text{M}} - \Omega_{\text{B}} = 0.265 \pm 0.011$$

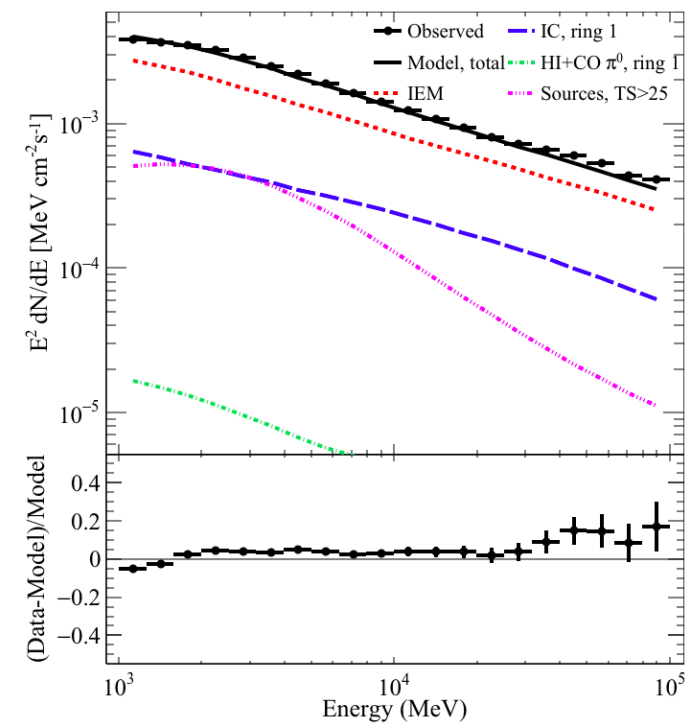
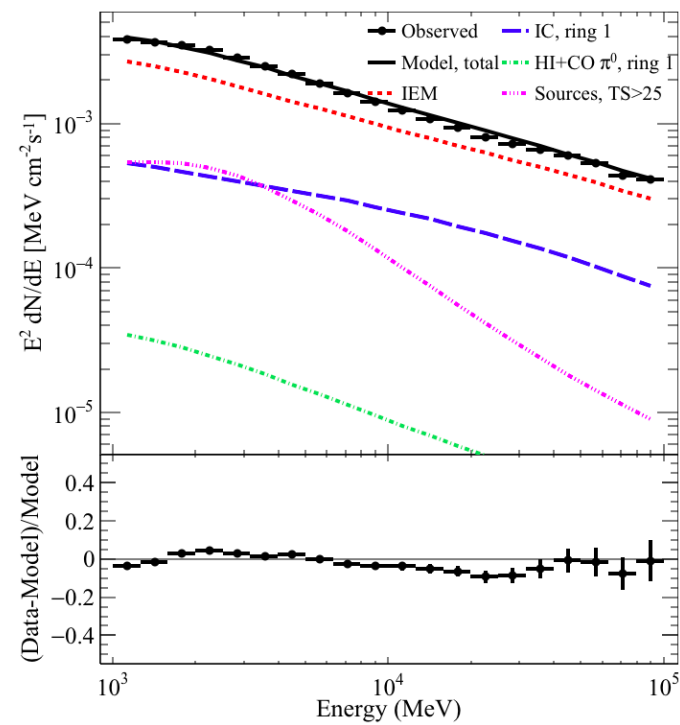
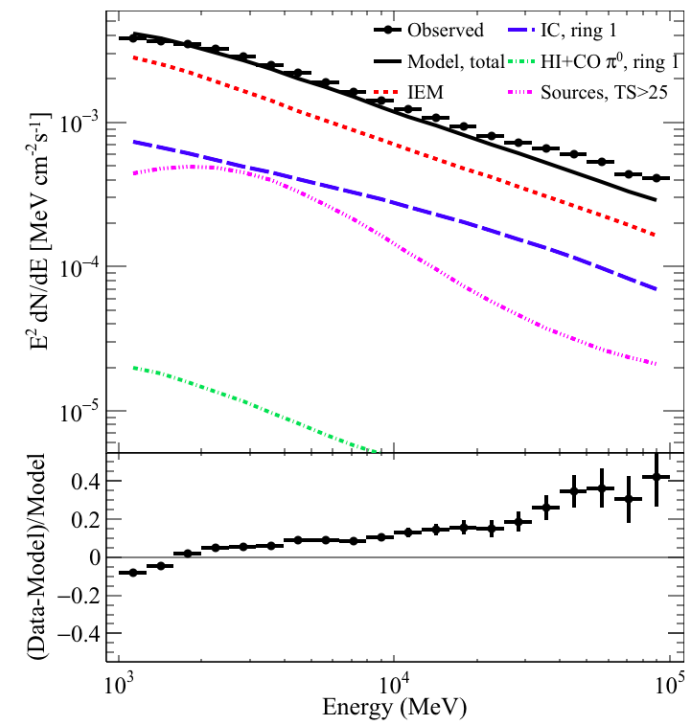
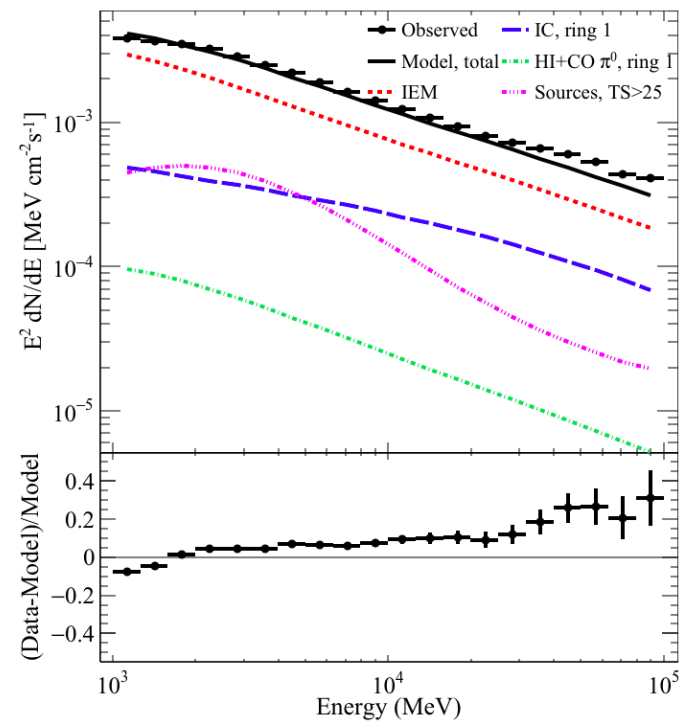
$$\Omega_{\chi} h^2 = 0.112 \pm 0.012$$



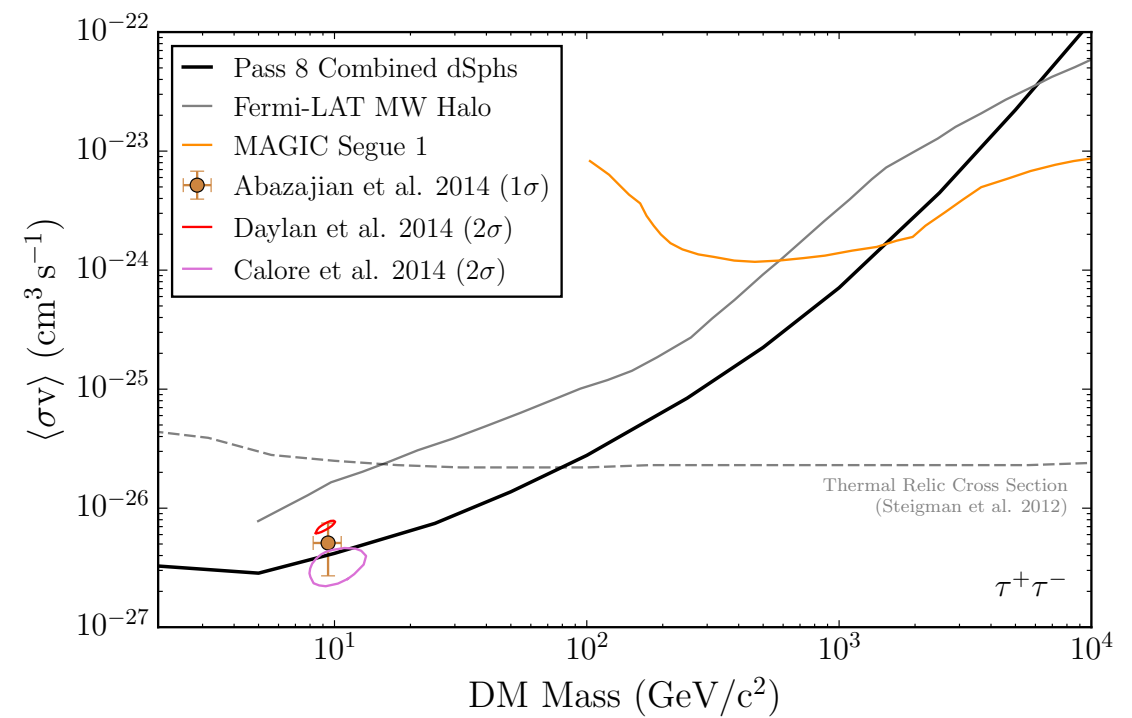
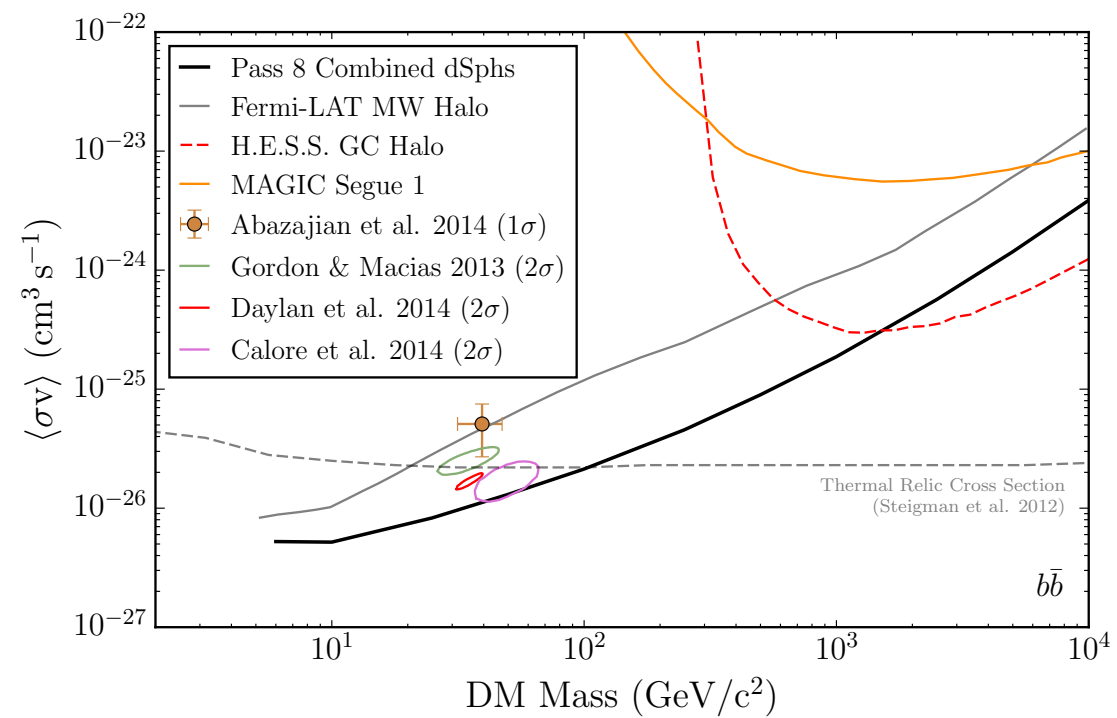
## DM/ Direct searches



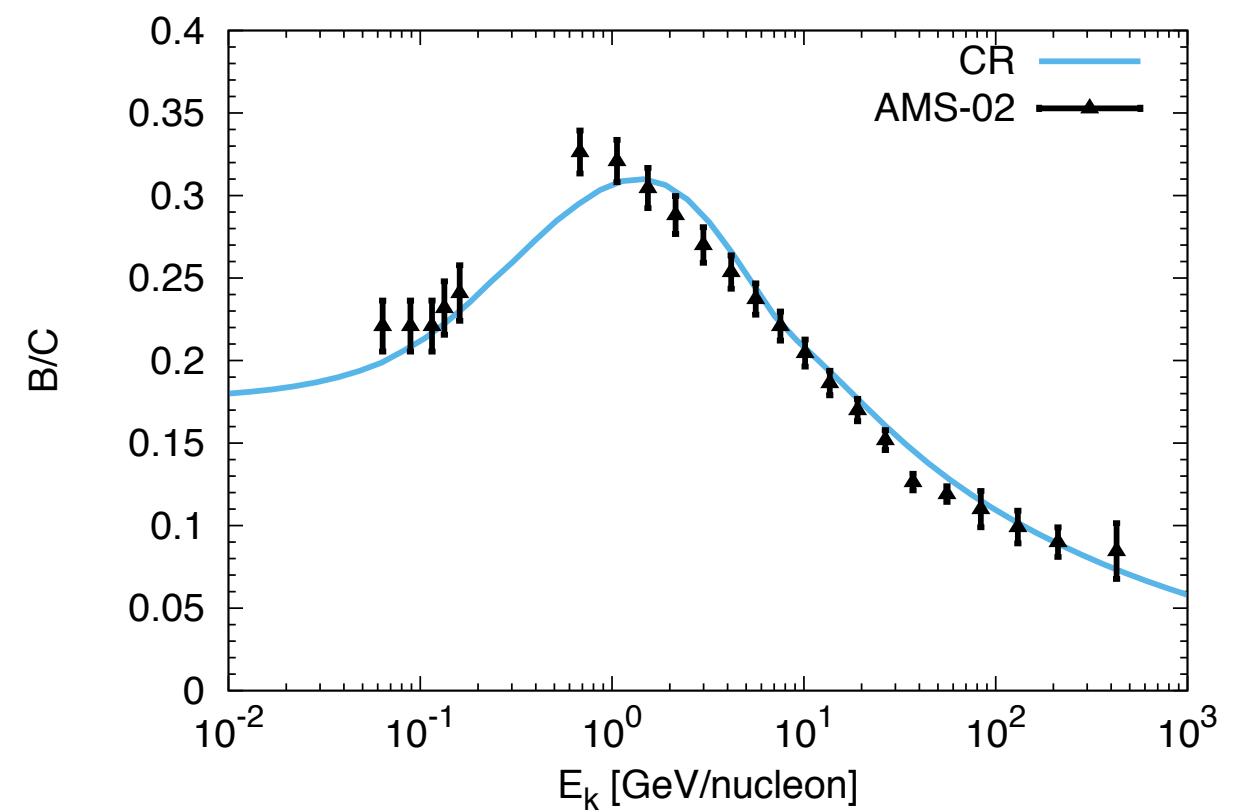
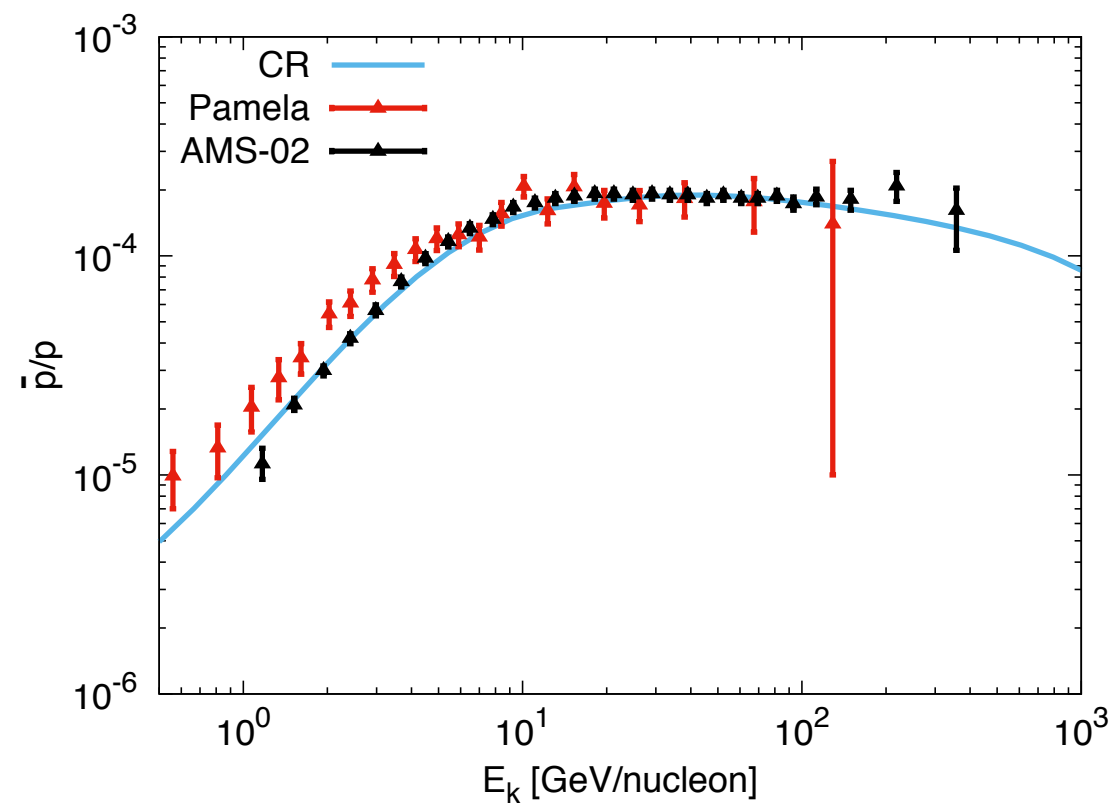
# DM/ Indirect searches gammas GC



# DM/ Indirect searches gammas dSph



# DM/ Indirect searches antiprotons



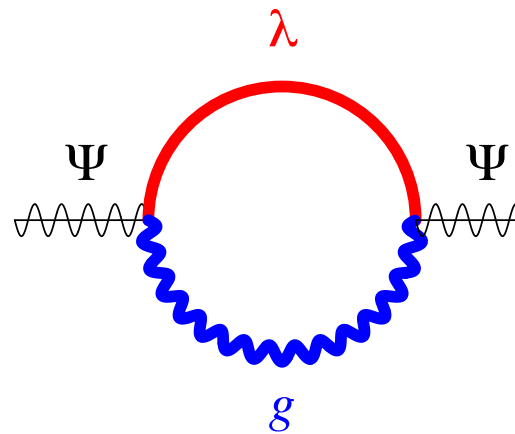
# DM searches status

- **Direct searches, up to date, have found nothing**
- **Indirect searches uncertainties: halo, direct x-section etc**
- **Neutralino models “currently not favoured” by LHC SUSY searches**
- **— Gravitino**

# Gravitino Dark Matter

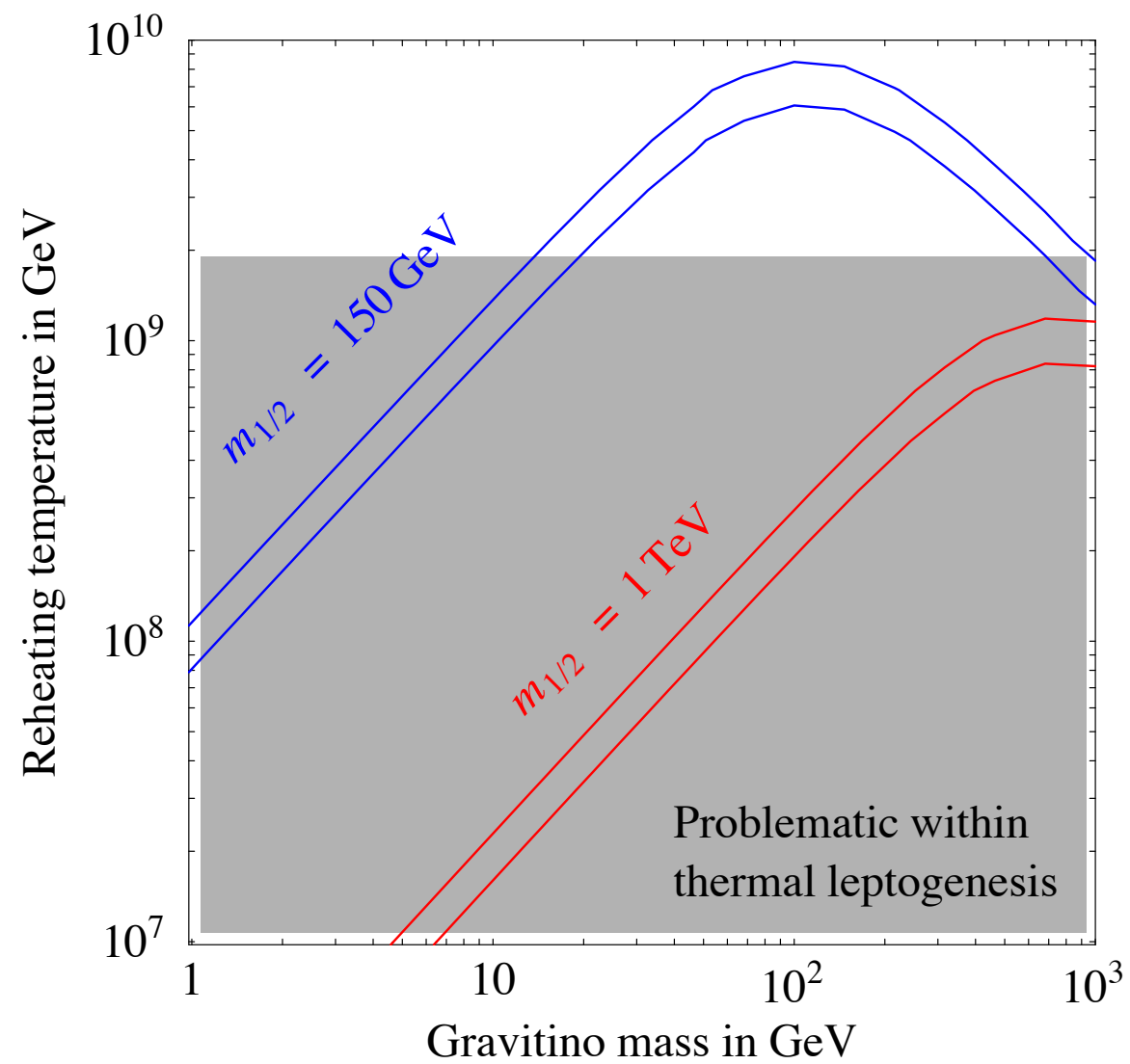
## Gravitino thermal production

$$Y_{3/2}(T) \simeq 4.48 \times 10^{-11} \left( \frac{T_{\text{reh}}}{10^{10} \text{ GeV}} \right) \times \left[ \left( 1 + 0.558 \frac{m_{1/2}^2}{m_{3/2}^2} \right) - 0.011 \left( 1 + 3.062 \frac{m_{1/2}^2}{m_{3/2}^2} \right) \ln \left( \frac{T_{\text{reh}}}{10^{10} \text{ GeV}} \right) \right]$$



**Important constraints from BBN data**

# Gravitino thermal abundance

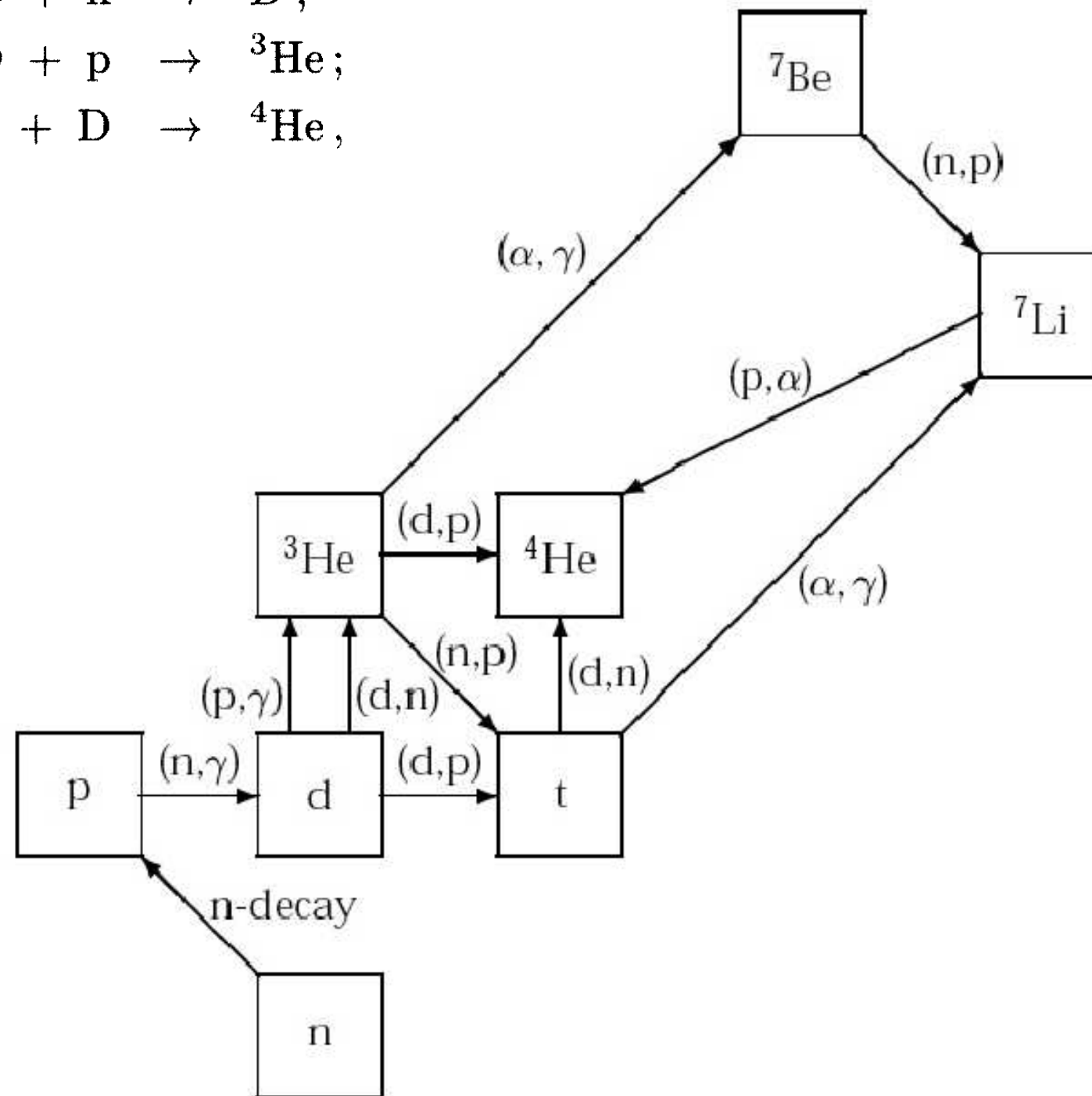
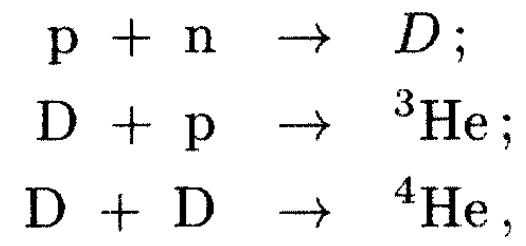


# Big Bang Nucleosynthesis

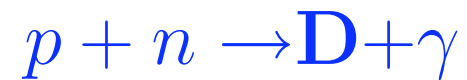
## Light Elements observed abundances:

- \*  $^4\text{He}$  observed in extragalactic HII regions:  
abundance by mass  $\sim 25\%$
- \*  $^7\text{Li}$  observed in the atmosphere of dwarf halo stars:  
abundance by number  $\sim 10^{-10}$
- \*  $\text{D}$  in quasars absorption systems (and locally):  
abundance by number  $\sim 3 \times 10^{-5}$
- \*  $^3\text{He}$  observed in solar wind, meteorites, and in ISM:  
abundance by number  $\sim 10^{-5}$





# Nucleosynthesis Delayed (Deuterium Bottleneck)



$$\Gamma_p \sim n_B \sigma$$



$$\Gamma_d \sim n_\gamma \sigma e^{-E_B/T}$$

Nucleosynthesis begins when  $\Gamma_p \sim \Gamma_d$

$$\frac{n_\gamma}{n_B} e^{-E_B/T} \sim 1 \quad @ \ T \sim 0.1 \text{ MeV}$$

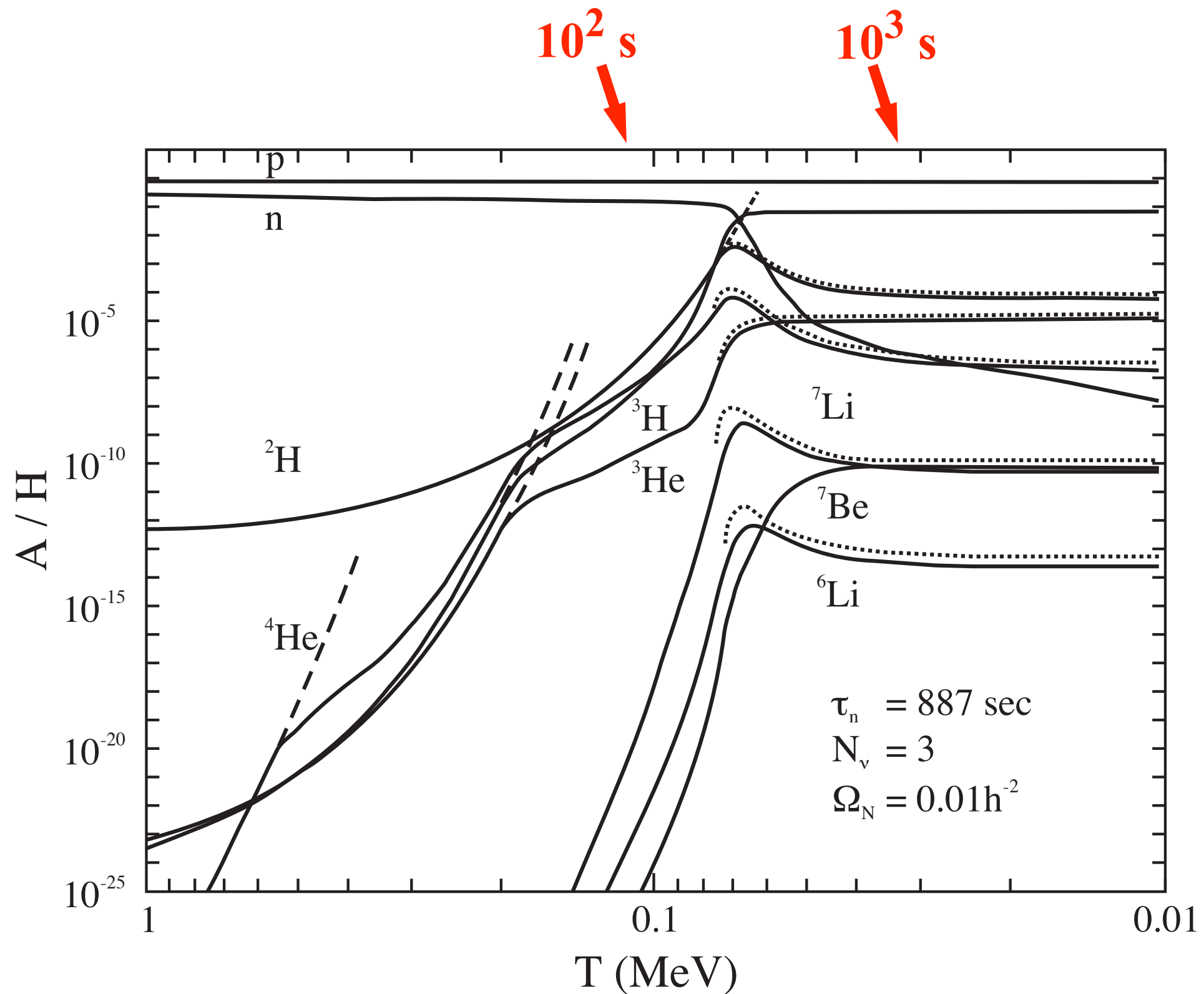
All neutrons  $\rightarrow$   $^4\text{He}$

$$Y_p = \frac{2(n/p)}{1 + (n/p)} \simeq 25\%$$

Remainder:

$\mathbf{D}$ ,  $^3\text{He} \sim 10^{-5}$  and  $^7\text{Li} \sim 10^{-10}$  by number

# *Light elements production*



GDM

## Gravitino field representation

The four polarisation states of the gravitino field in the momentum space in terms of spin-1 and spin-1/2 components:

$$\begin{aligned}\psi_\mu(\mathbf{k}, \frac{3}{2}) &= u(\mathbf{k}, \frac{1}{2}) \epsilon_\mu(\mathbf{k}, 1), \\ \psi_\mu(\mathbf{k}, \frac{1}{2}) &= \sqrt{\frac{1}{3}} u(\mathbf{k}, -\frac{1}{2}) \epsilon_\mu(\mathbf{k}, 1) + \sqrt{\frac{2}{3}} u(\mathbf{k}, \frac{1}{2}) \epsilon_\mu(\mathbf{k}, 0) \\ \psi_\mu(\mathbf{k}, -\frac{1}{2}) &= \sqrt{\frac{1}{3}} u(\mathbf{k}, \frac{1}{2}) \epsilon_\mu(\mathbf{k}, -1) + \sqrt{\frac{2}{3}} u(\mathbf{k}, -\frac{1}{2}) \epsilon_\mu(\mathbf{k}, 0) \\ \psi_\mu(\mathbf{k}, -\frac{3}{2}) &= u(\mathbf{k}, -\frac{1}{2}) \epsilon_\mu(\mathbf{k}, -1),\end{aligned}$$

From the field equation for the spin-3/2 particle, the so called [Rarita-Schwinger equation](#), we get three equations in momentum space, which we have to fulfill by the right choice of the spin-1 and spin-1/2 component representations

$$\begin{aligned}\gamma^\mu \psi_\mu(\mathbf{k}, \lambda) &= 0 \\ k^\mu \psi_\mu(\mathbf{k}, \lambda) &= 0 \\ (\not{k} - m_{\tilde{G}}) \psi_\mu(\mathbf{k}, \lambda) &= 0\end{aligned}$$

GDM

## Gravitino Interactions with the MSSM

The relevant supergravity Lagrangian reads

$$\mathcal{L}_{\tilde{G}, \text{int}}^{(\alpha)} = -\frac{i}{\sqrt{2}M_{\text{P}}} \left[ \mathcal{D}_{\mu}^{(\alpha)} \phi^{*i} \tilde{\bar{G}}_{\nu} \gamma^{\mu} \gamma^{\nu} \chi_L^i - \mathcal{D}_{\mu}^{(\alpha)} \phi^i \bar{\chi}_L^i \gamma^{\nu} \gamma^{\mu} \tilde{G}_{\nu} \right] \\ - \frac{i}{8M_{\text{P}}} \tilde{\bar{G}}_{\mu} [\gamma^{\rho}, \gamma^{\sigma}] \gamma^{\mu} \lambda^{(\alpha) a} F_{\rho\sigma}^{(\alpha) a} ,$$

with the covariant derivative given by

$$\mathcal{D}_{\mu}^{(\alpha)} \phi^i = \partial_{\mu} \phi^i + ig_{\alpha} A_{\mu}^{(\alpha) a} T_{a, ij}^{(\alpha)} \phi^j ,$$

and the field strength tensor  $F_{\mu\nu}^{(\alpha) a}$  reads

$$F_{\mu\nu}^{(\alpha) a} = \partial_{\mu} A_{\nu}^{(\alpha) a} - \partial_{\nu} A_{\mu}^{(\alpha) a} - g_{\alpha} f^{(\alpha) abc} A_{\mu}^{(\alpha) b} A_{\nu}^{(\alpha) c} .$$

The index  $\alpha$  corresponds to the three groups  $\text{U}(1)_Y$ ,  $\text{SU}(2)_I$ , and  $\text{SU}(3)_c$  with  $a = 1, 3, 8$  and  $i = 1, 2, 3$ , respectively.

# Steps

- \* Calculate the partial and the total decay widths
- \* Employ **PYTHIA** event generator to simulate the **EM** and **HD** products of Z, Higgs bosons, quarks and taus
- \* Incorporate in the **BBN** code the effects of the **EM** and **HD** injections
- \* Estimate for **each point of the SUSY parameter space** the light element abundances
- \* Important point: Bound state phenomena of charged metastable particles like stau!

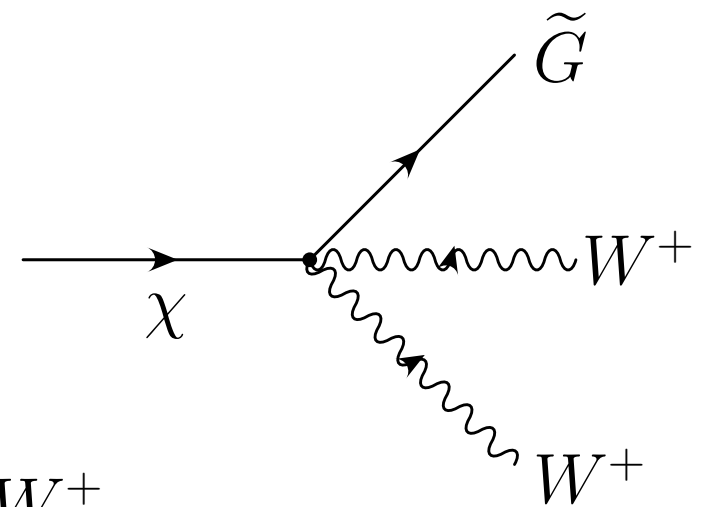
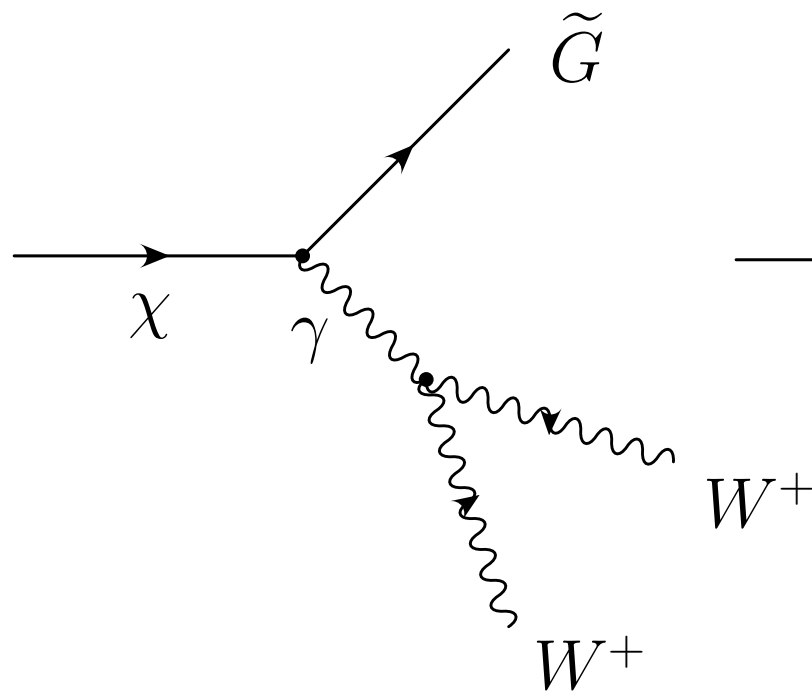
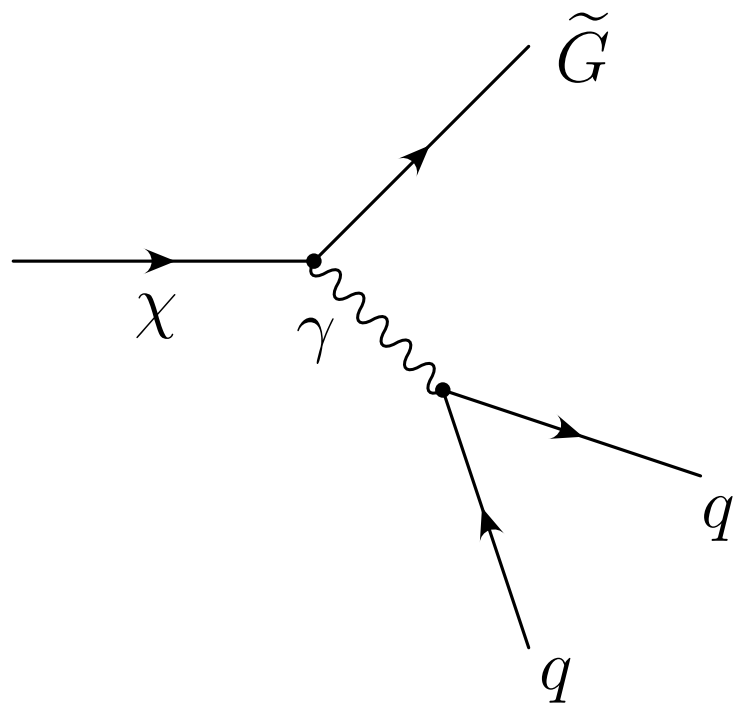
# NSP decays

$\chi$  NSP

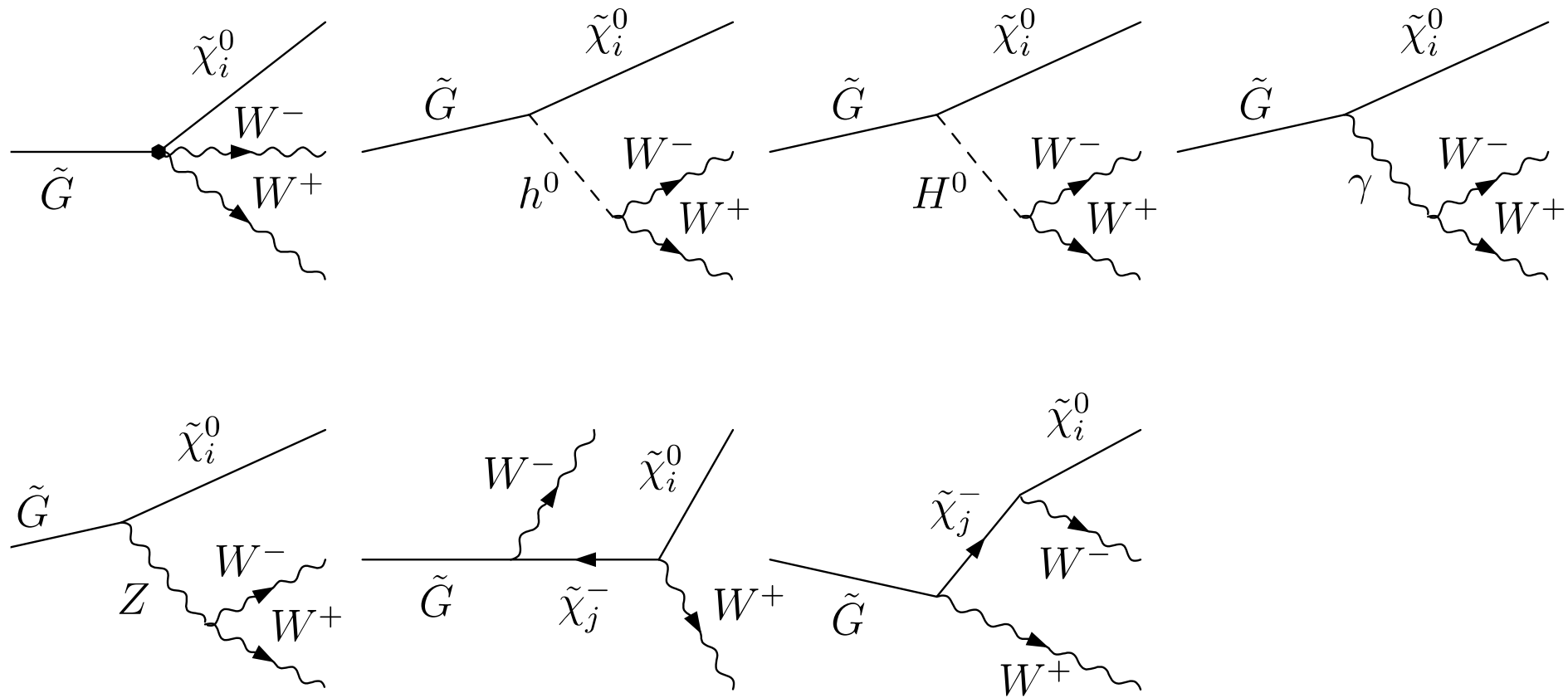
$$\chi \rightarrow \tilde{G} \gamma$$

$$\chi \rightarrow \tilde{G} Z$$

$$\chi \rightarrow \tilde{G} H_i$$



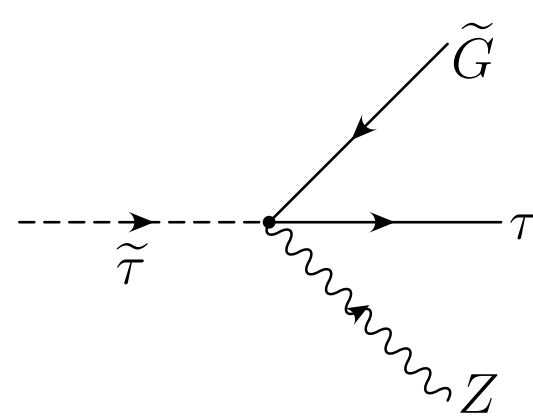
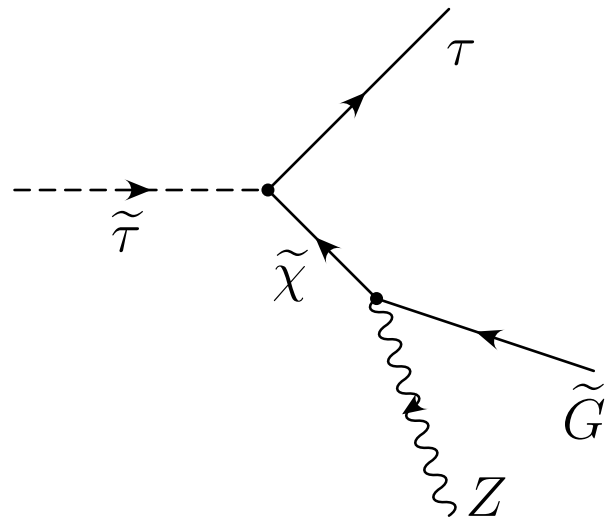
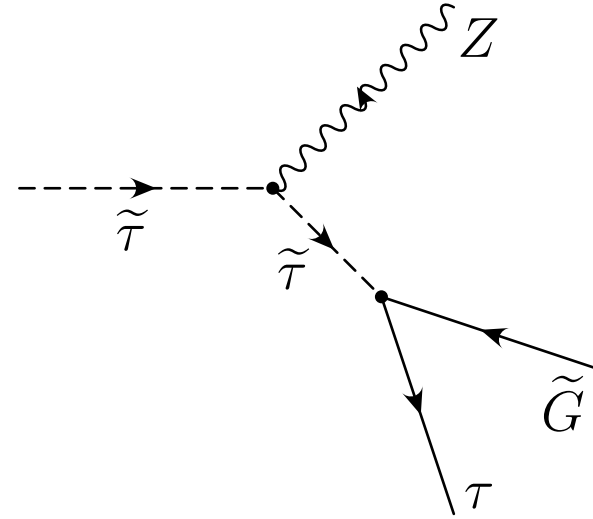
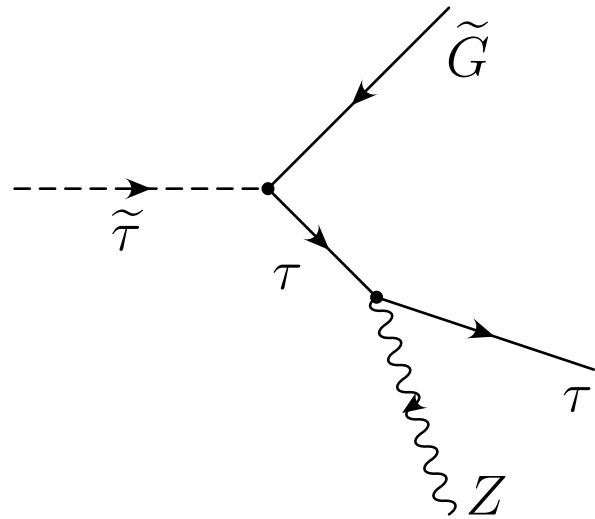
$$\tilde{G} \rightarrow \tilde{\chi}_i^0 W^- W^+$$





# $\tilde{\tau}$ NSP

$$\tilde{\tau} \rightarrow \tilde{G} \tau$$

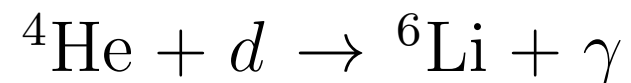


process $\tilde{\chi}_1^0 \rightarrow \tilde{G}XY$	number of graphs	first decay $\tilde{\chi}_1^0 \rightarrow \tilde{X}Y$	possible resonances
$\tilde{G}f\bar{f}$	7	$\tilde{G}(h^0, H^0, A^0, \gamma, Z^0), f\tilde{f}_l^*, \tilde{f}_l\bar{f}$	$h^0, H^0, A^0, Z^0$
$\tilde{G}Z^0Z^0$	4	$\tilde{G}(h^0, H^0), Z^0\tilde{\chi}_k^0, \tilde{\chi}_k^0Z^0$	$H^0$
$\tilde{G}Z^0\gamma$	1	$\tilde{\chi}_k^0Z^0$	—
$\tilde{G}W^+W^-$	$6 + 4_{pt}$	$\tilde{G}(h^0, H^0, \gamma, Z^0), W^+\tilde{\chi}_j^-, \tilde{\chi}_j^+W^-$	$H^0$
$\tilde{G}Z^0h^0$	$4 + 4_{pt}$	$\tilde{G}(A^0, Z^0), Z^0\tilde{\chi}_k^0, \tilde{\chi}_k^0h^0$	$A^0$
$\tilde{G}Z^0H^0$	$4 + 4_{pt}$	$\tilde{G}(A^0, Z^0), Z^0\tilde{\chi}_k^0, \tilde{\chi}_k^0H^0$	$A^0$
$\tilde{G}Z^0A^0$	$4 + 4_{pt}$	$\tilde{G}(h^0, H^0), \tilde{\chi}_k^0Z^0, A^0\tilde{\chi}_k^0$	$H^0$
$\tilde{G}\gamma h^0$	1	$\tilde{\chi}_k^0h^0$	—
$\tilde{G}\gamma H^0$	1	$\tilde{\chi}_k^0H^0$	—
$\tilde{G}\gamma A^0$	1	$\tilde{\chi}_k^0A^0$	—
$\tilde{G}W^+H^-$	$5 + 4_{pt}$	$\tilde{G}(h^0, H^0, A^0), W^+\tilde{\chi}_j^-, \tilde{\chi}_j^+H^-$	$H^0, A^0$
$\tilde{G}W^-H^+$	$5 + 4_{pt}$	$\tilde{G}(h^0, H^0, A^0), W^-\tilde{\chi}_j^+, \tilde{\chi}_j^-H^+$	$H^0, A^0$
$\tilde{G}h^0h^0$	4	$\tilde{G}(h^0, H^0), h^0\tilde{\chi}_k^0, \tilde{\chi}_k^0h^0$	$H^0$
$\tilde{G}H^0H^0$	4	$\tilde{G}(h^0, H^0), H^0\tilde{\chi}_k^0, \tilde{\chi}_k^0H^0$	—
$\tilde{G}h^0H^0$	4	$\tilde{G}(h^0, H^0), h^0\tilde{\chi}_k^0, \tilde{\chi}_k^0H^0$	—
$\tilde{G}A^0A^0$	4	$\tilde{G}(h^0, H^0), A^0\tilde{\chi}_k^0, \tilde{\chi}_k^0A^0$	$H^0$
$\tilde{G}h^0A^0$	3	$\tilde{G}(A^0, Z^0), h^0\tilde{\chi}_k^0$	—
$\tilde{G}H^0A^0$	3	$\tilde{G}(A^0, Z^0), H^0\tilde{\chi}_k^0$	—
$\tilde{G}H^+H^-$	6	$\tilde{G}(h^0, H^0, \gamma, Z^0), H^+\tilde{\chi}_j^-, \tilde{\chi}_j^+H^-$	$H^0$

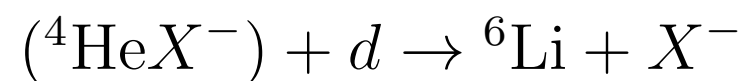
19 individuel  
decay channels

# Bound-state catalysis

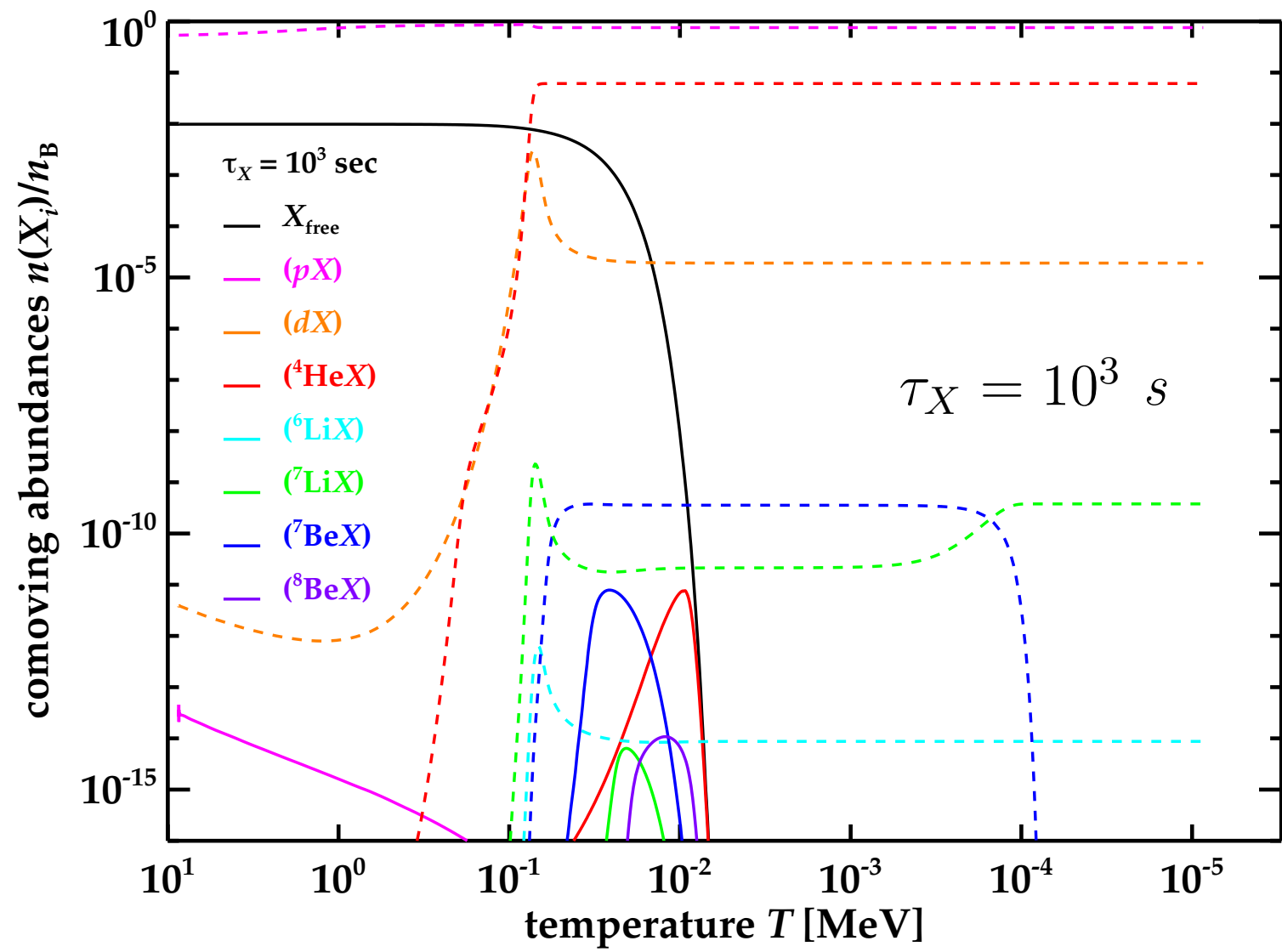
**standard BBN reaction is suppressed**



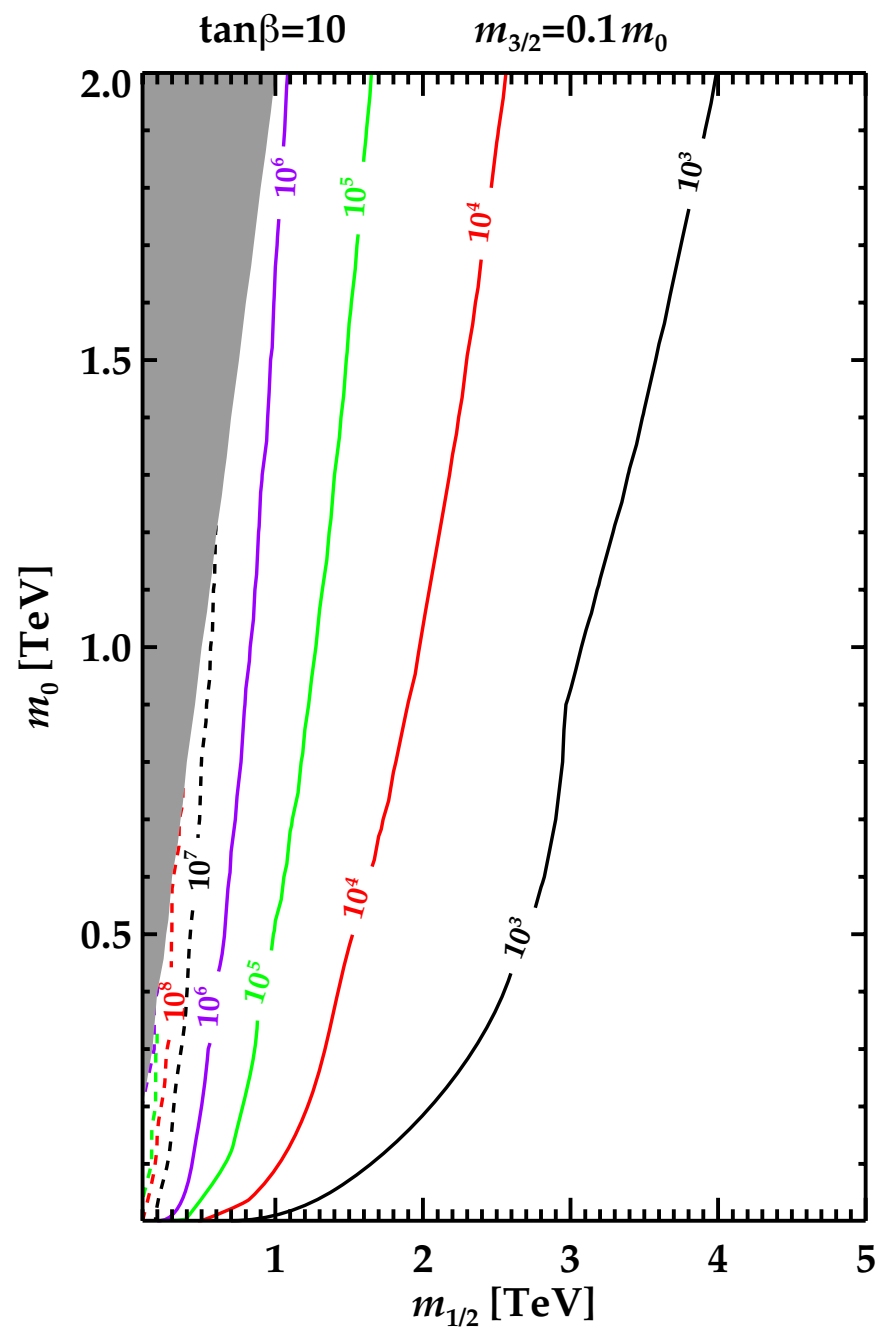
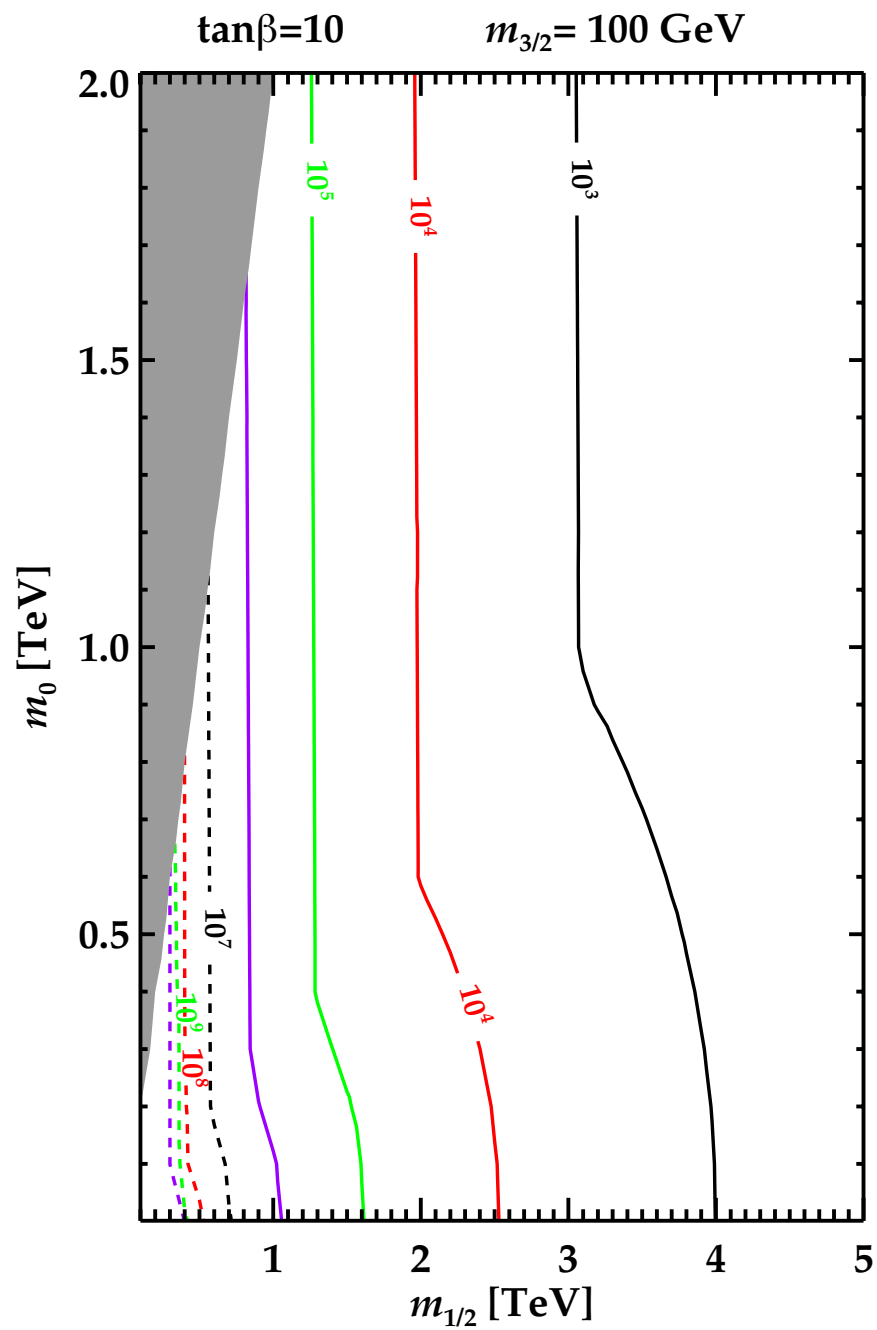
**the bound state reaction is not suppressed by photon emission**



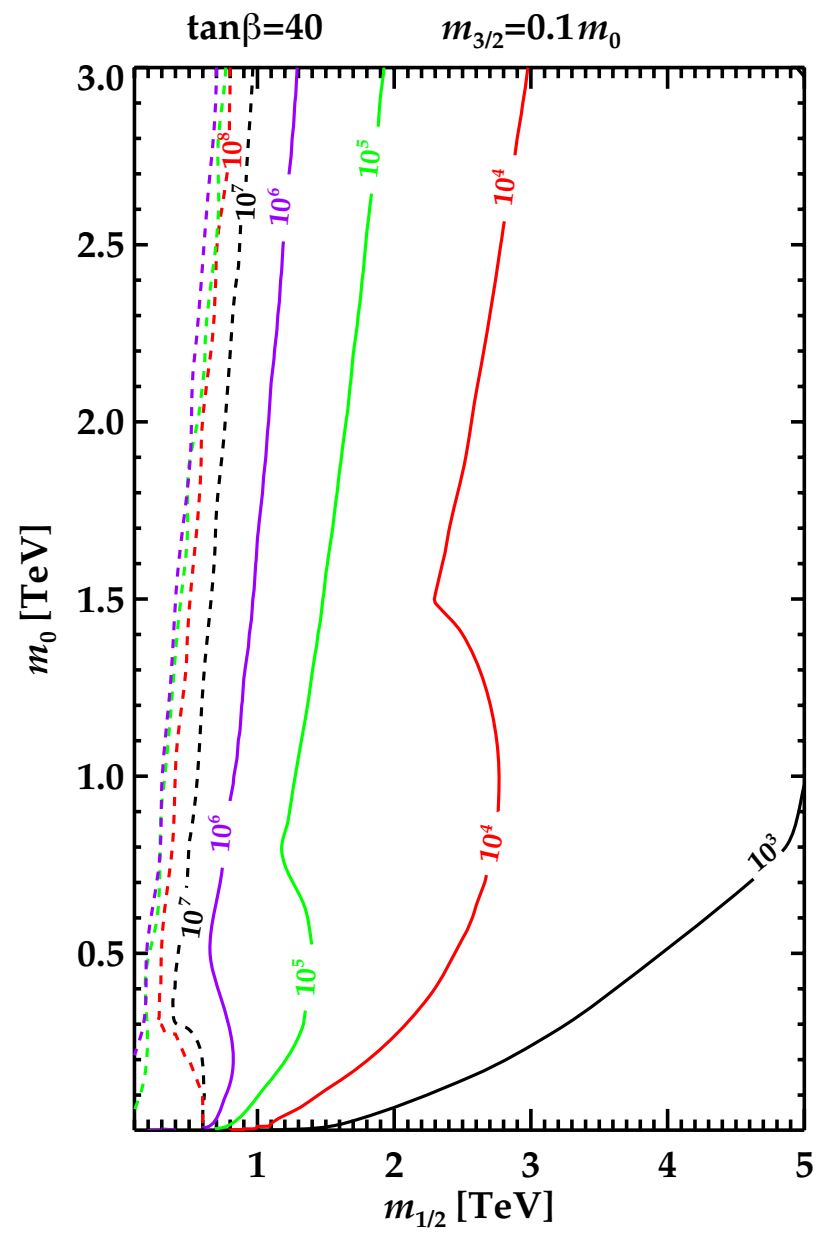
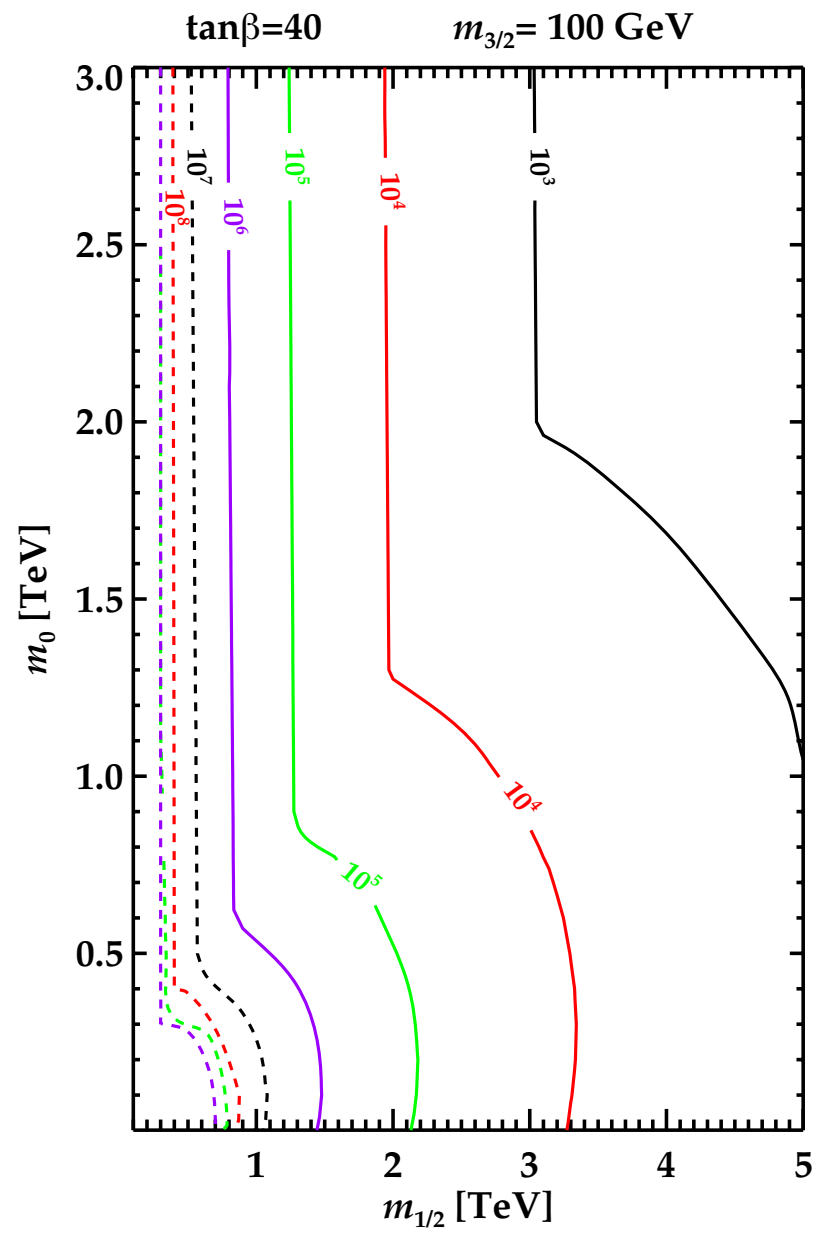
**possible solution of Li puzzle**



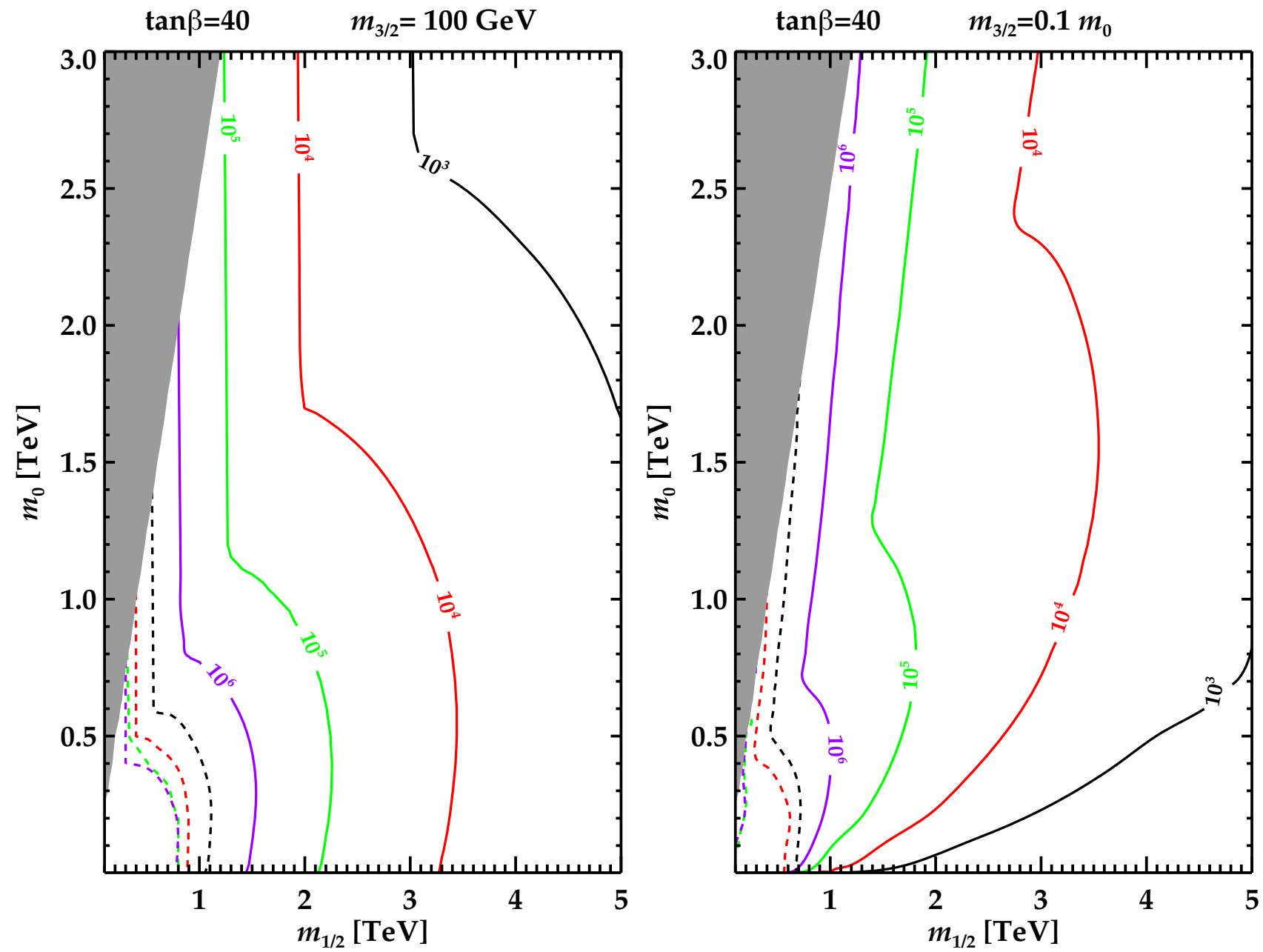
$$A_0 = 2.5 m_0$$

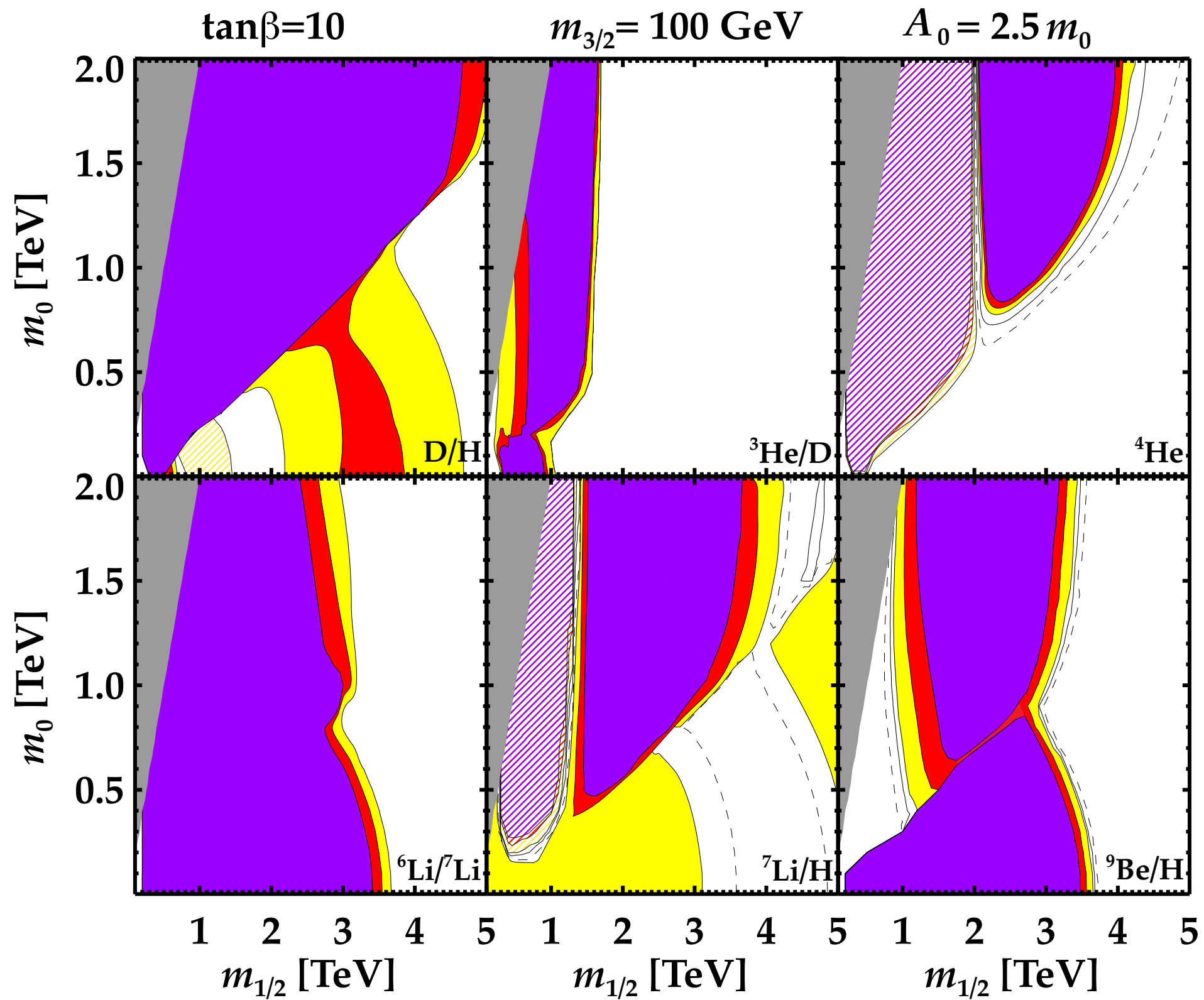


$$A_0 = 2 m_0.$$



$$A_0 = 2.5 m_0$$

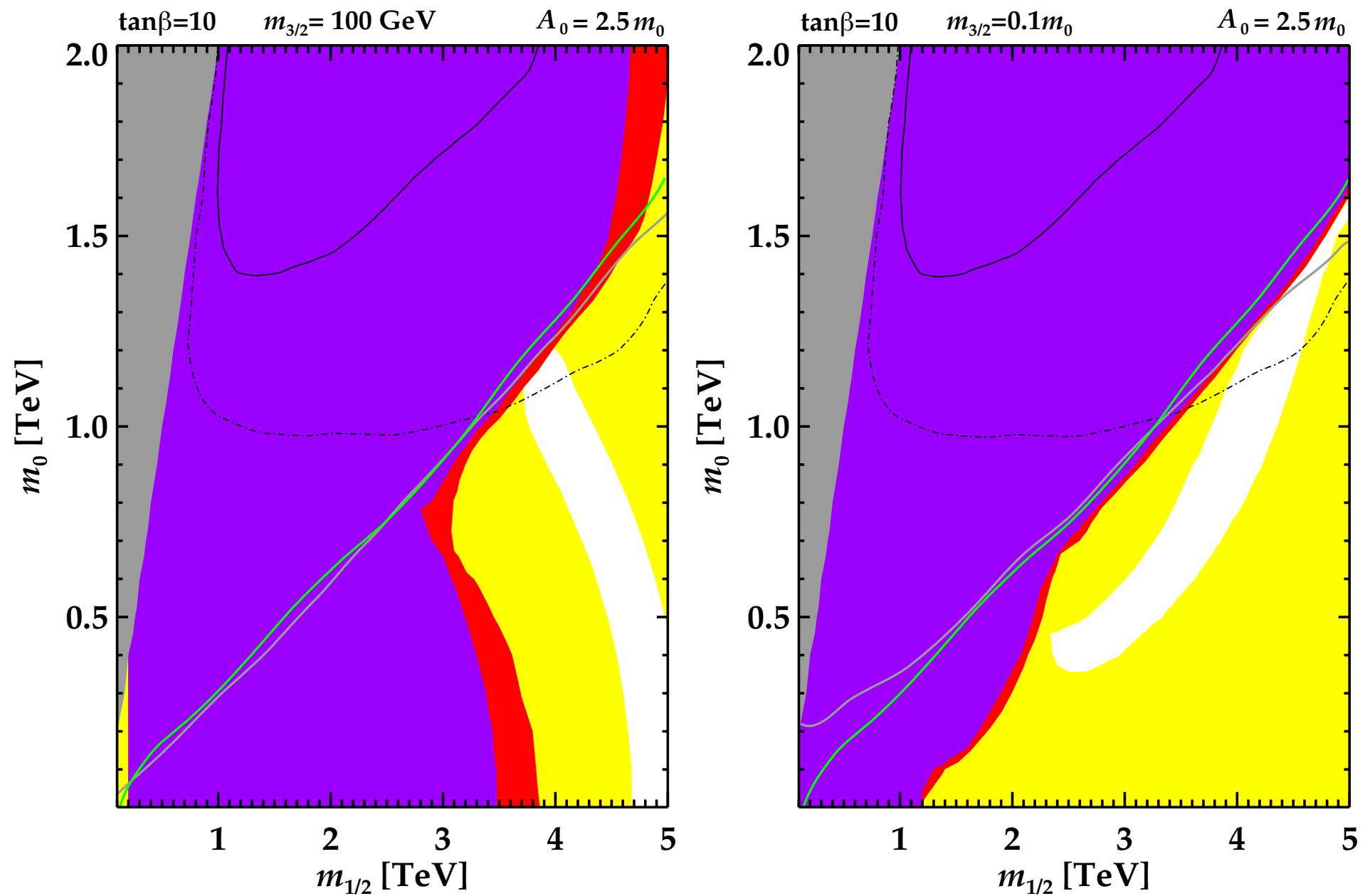




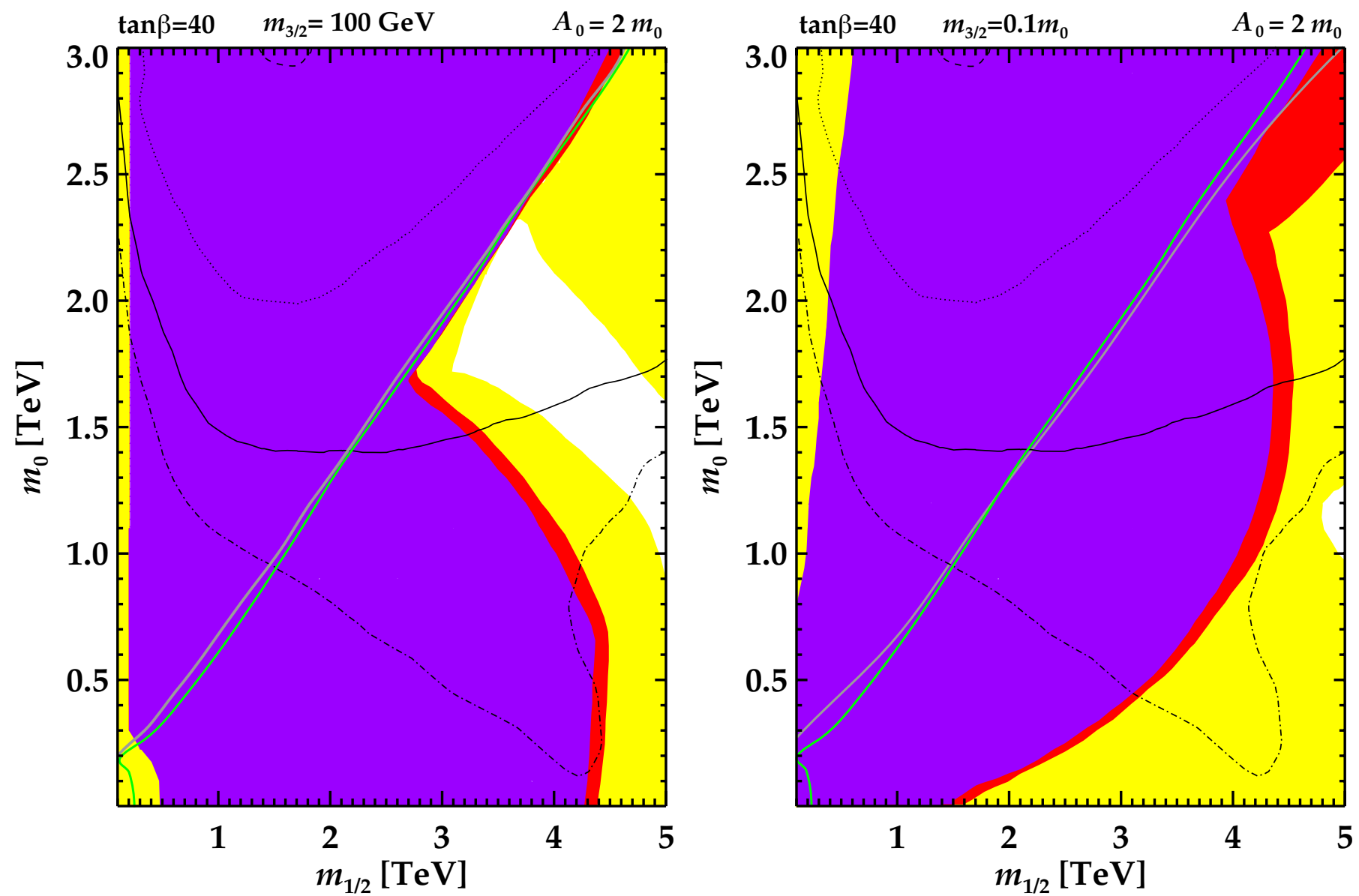


GDM/ results

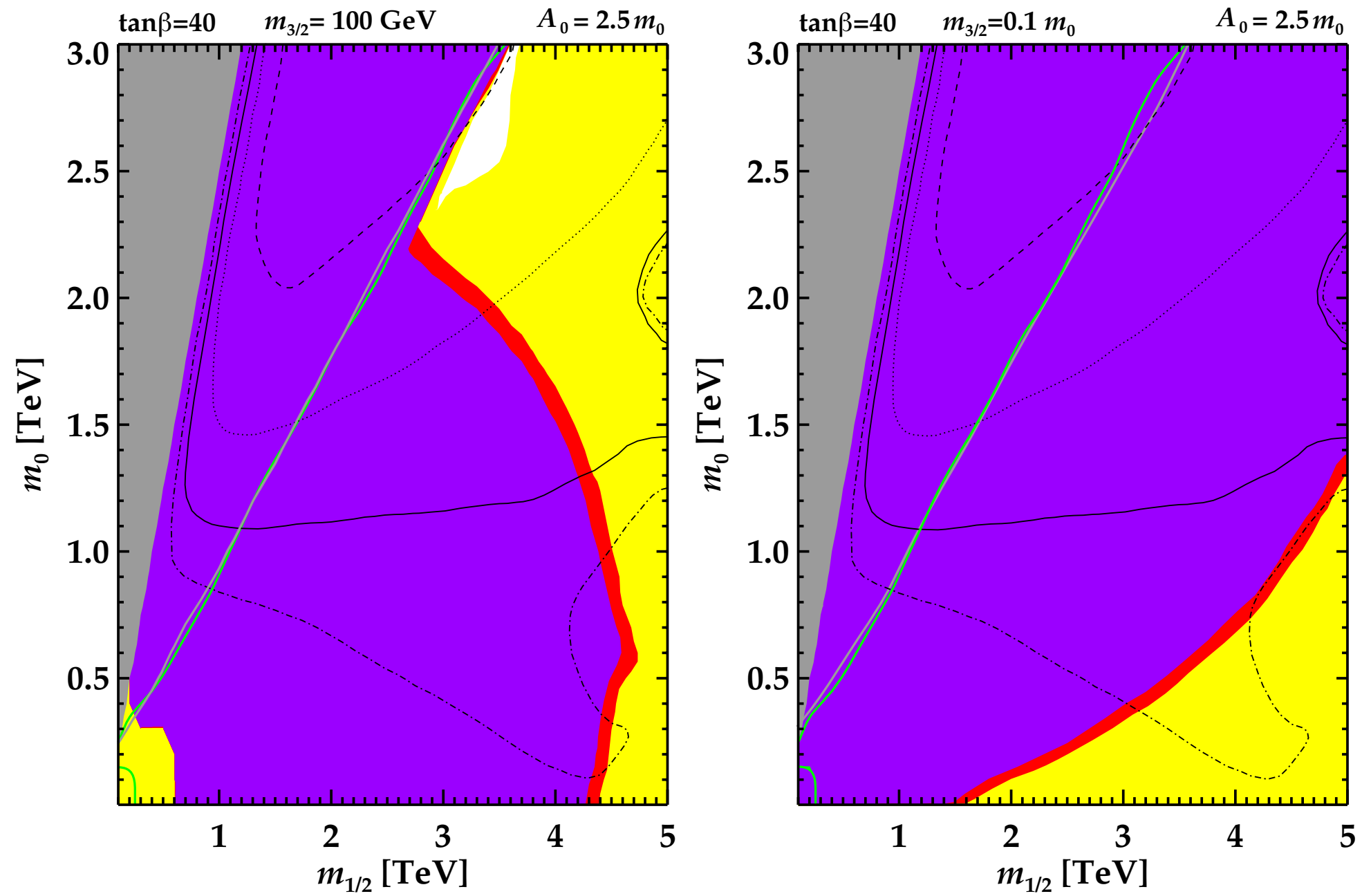
# Gravitino DM parameter space

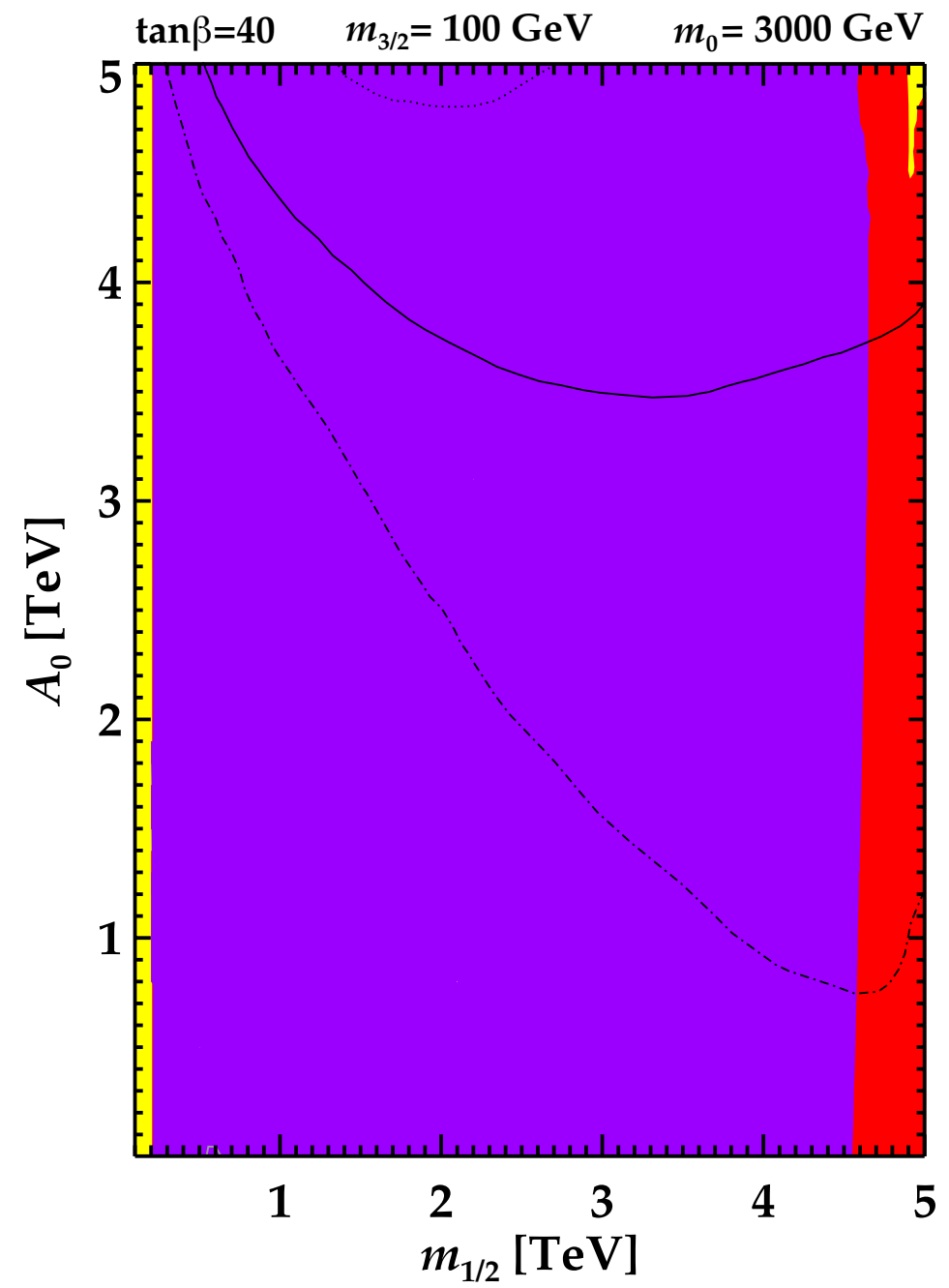
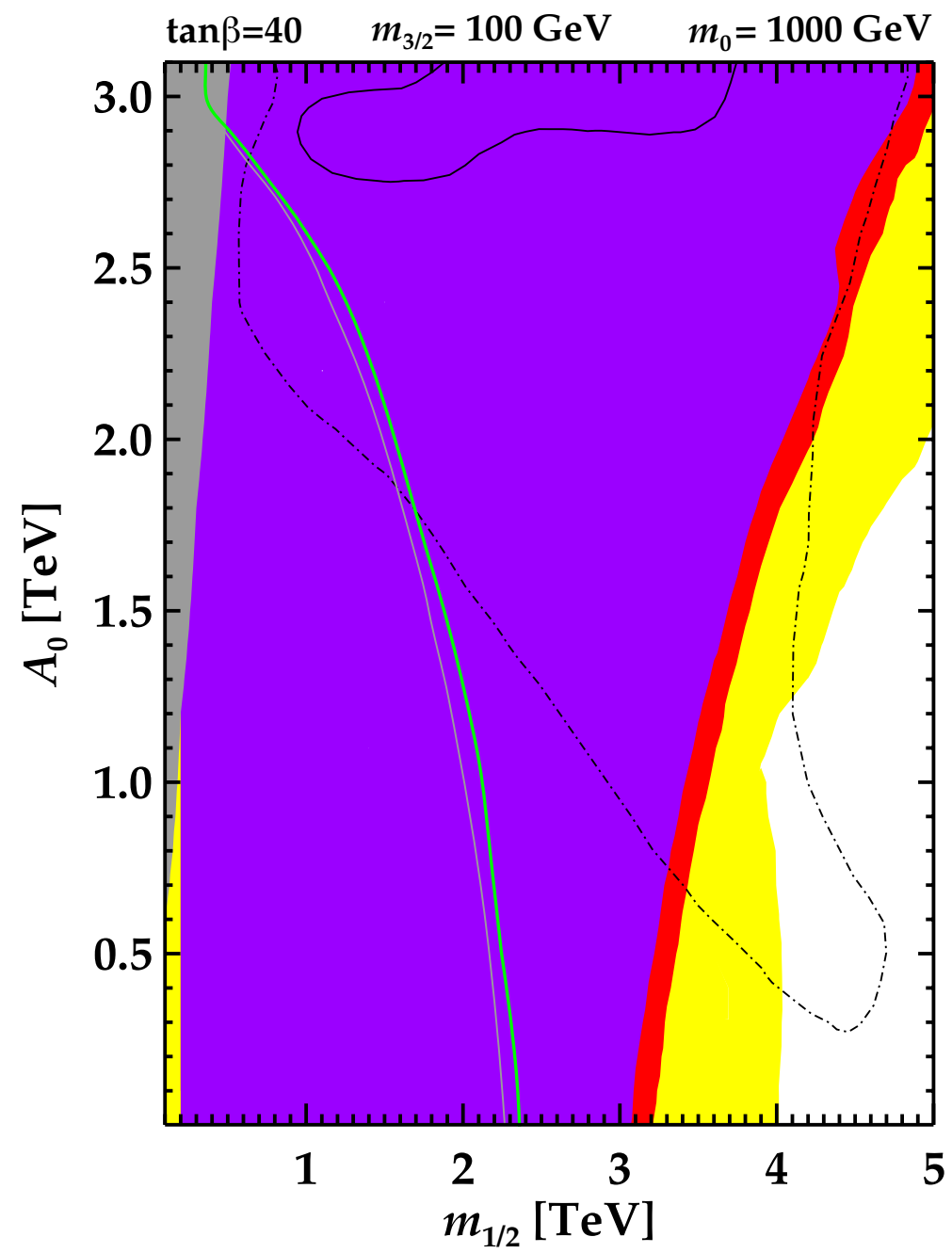


# GDM/ results



# GDM/ results





# Summary & Prospects

- \* Gravitino is natural candidate in SUGRA models for DM
- \* GDM scenario can be quite involved (bound state effects)
- \* Inclusion of Gravitino thermal production is important to set bound to reheating temperature
- \* DM scenario and phenomenology can open an window to Planck scale
- \* **Many-Many Thanks and Congratulations to Keith on his 60th birthday !!**